



# CERFACS

CENTRE EUROPÉEN DE RECHERCHE ET DE FORMATION AVANCÉE EN **CALCUL SCIENTIFIQUE**

## CFD tools specificities for H2 flames and available/recommended models in AVBP for hydrogen/air flames.

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H2 Week • IMFT • 2024.02.27



# Chemistry



# Chemical Kinetics

*Just this once, H<sub>2</sub> is simpler to model than conventional fuels*

## Conventional fuels (e.g. kerosene)

- Global chemistries:

$$N_{\text{spec}} \simeq 6, N_{\text{reac}} \simeq 2$$

- Analytically Reduced Chemistries

$$N_{\text{spec}} \simeq 20 - 30, N_{\text{reac}} \simeq 200 - 500$$

- Detailed chemistries (e.g. CRECK [1]):

$$N_{\text{spec}} \simeq 600, N_{\text{reac}} \simeq 30000$$

## Hydrogen

- Global chemistries:

$$N_{\text{spec}} \simeq 4, N_{\text{reac}} \simeq 1$$

- ~~Analytically Reduced Chemistries~~

- Detailed chemistries (e.g., UCSD aka San Diego [2]):

$$N_{\text{spec}} \simeq 9, N_{\text{reac}} \simeq 21$$

[1] Bieleveld et al., *Proc. Comb. Inst.*, 2009

[2] Saxena & Williams, *Combust. Flame*, 2006



# Transport properties



# Transport properties

## *Laminar viscosity of the gas*

- For conventional fuels, two dynamic viscosity laws are typically used:

### *Sutherland*

$$\mu = c_1 \frac{T^{3/2}}{T + c_2} \frac{T_{ref} + c_2}{T_{ref}^{3/2}}$$

### *Power law*

$$\mu = c_1 \left( \frac{T}{T_{ref}} \right)^b$$

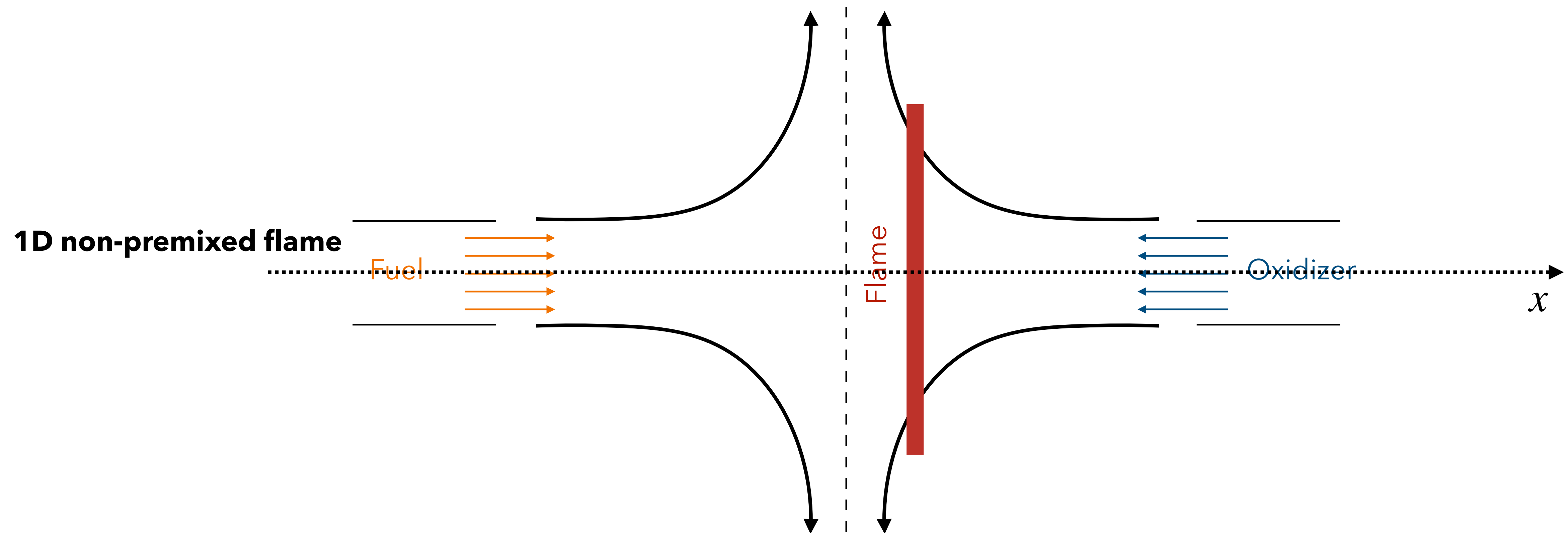
These laws are based on air.  
→ is this valid for H<sub>2</sub>/Air flames?



# Transport properties

## *Laminar viscosity of the gas*

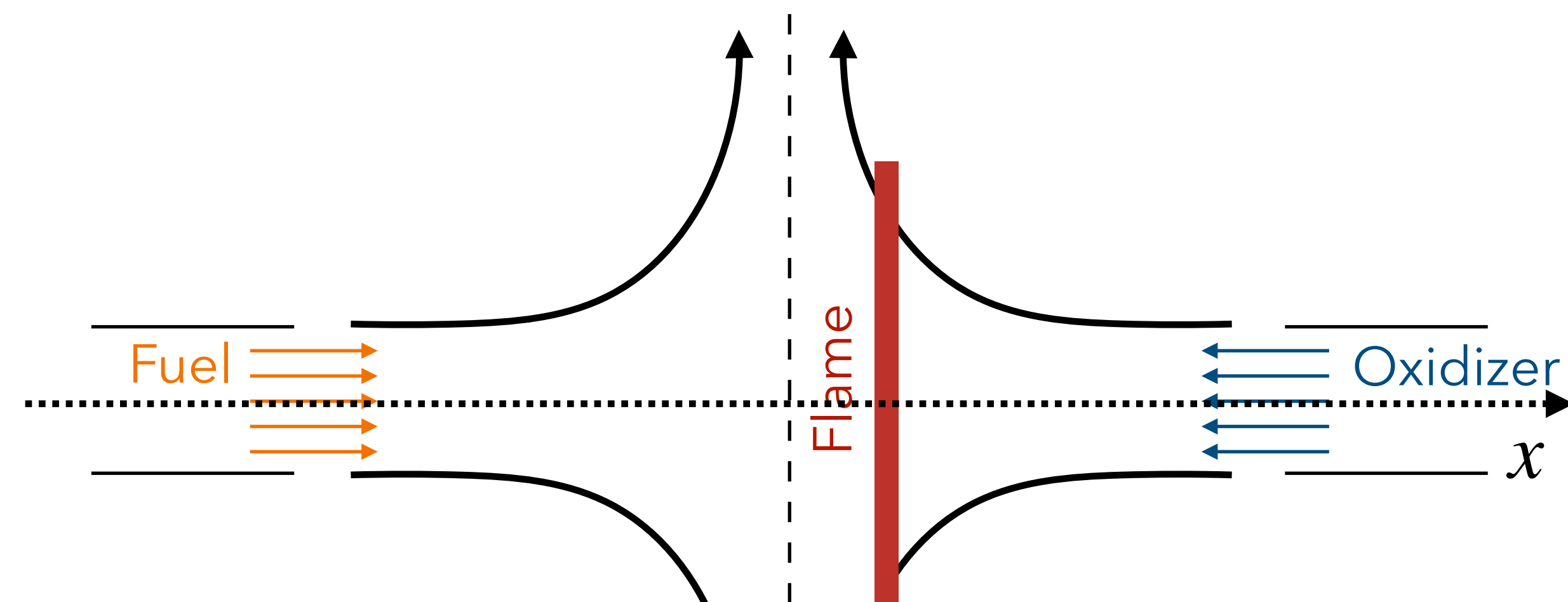
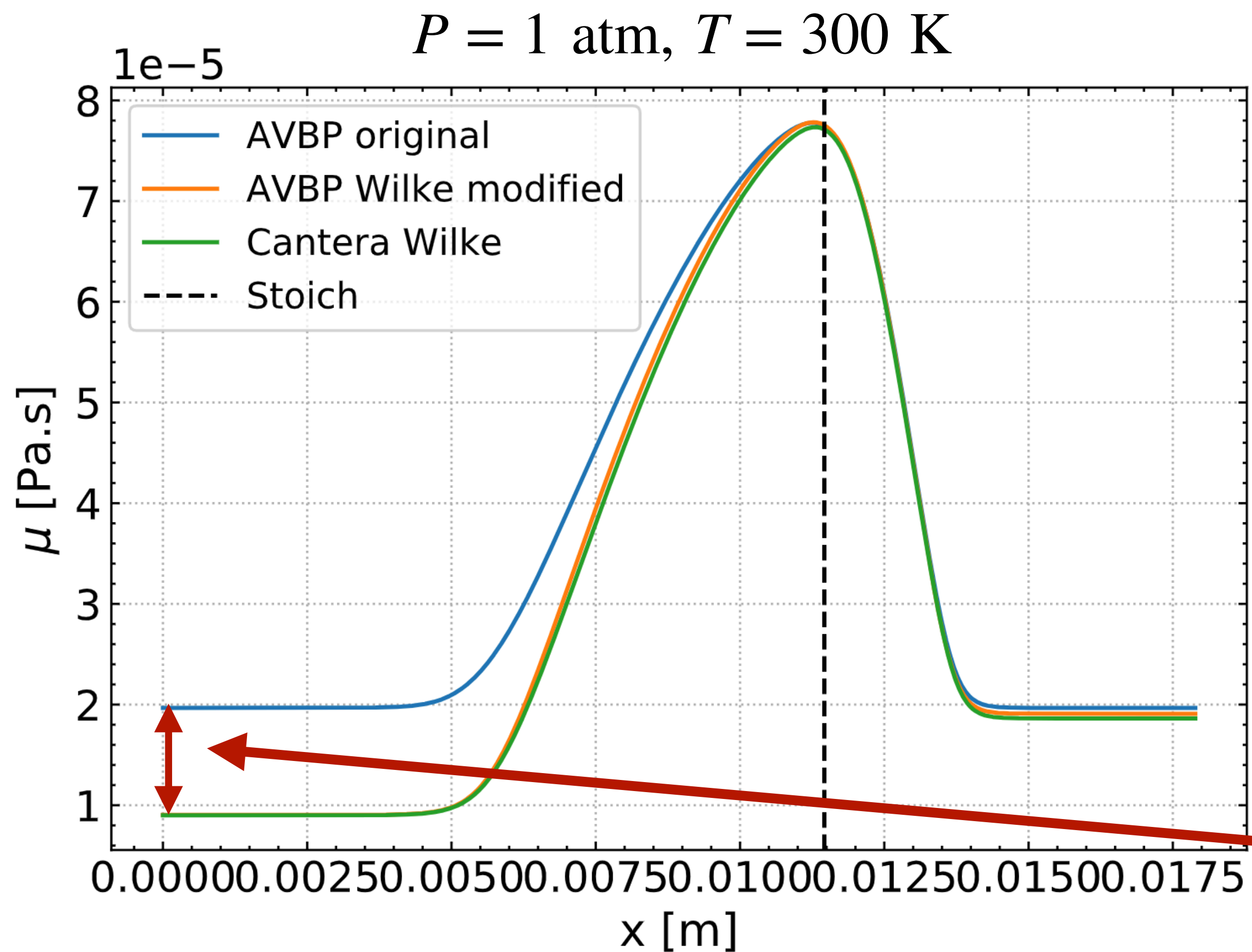
Let's consider a canonical counterflow non-premixed flame





# Transport properties

## Laminar viscosity of the gas



Laminar viscosity is wrong by a factor 2 in pure hydrogen



# Transport properties

## Laminar viscosity of the gas

Wilke's binary mixing between Pure H2 and Air

$$\mu = \frac{X_{H_2} \mu_{H_2}}{X_{H_2} + (1 - X_{H_2}) \phi_{H_2,air}} + \frac{(1 - X_{H_2}) \mu_{air}}{X_{H_2} \phi_{air,H_2} + 1 - X_{H_2}}$$

Power law

$$\phi_{air,H_2} = \frac{\left[ 1 + \left( \frac{\phi_{air}}{\phi_{H_2}} \right)^{1/2} \left( \frac{W_{H_2}}{W_{air}} \right)^{1/4} \right]^2}{2\sqrt{2} \left( 1 + \frac{W_{air}}{W_{H_2}} \right)^{1/2}}$$

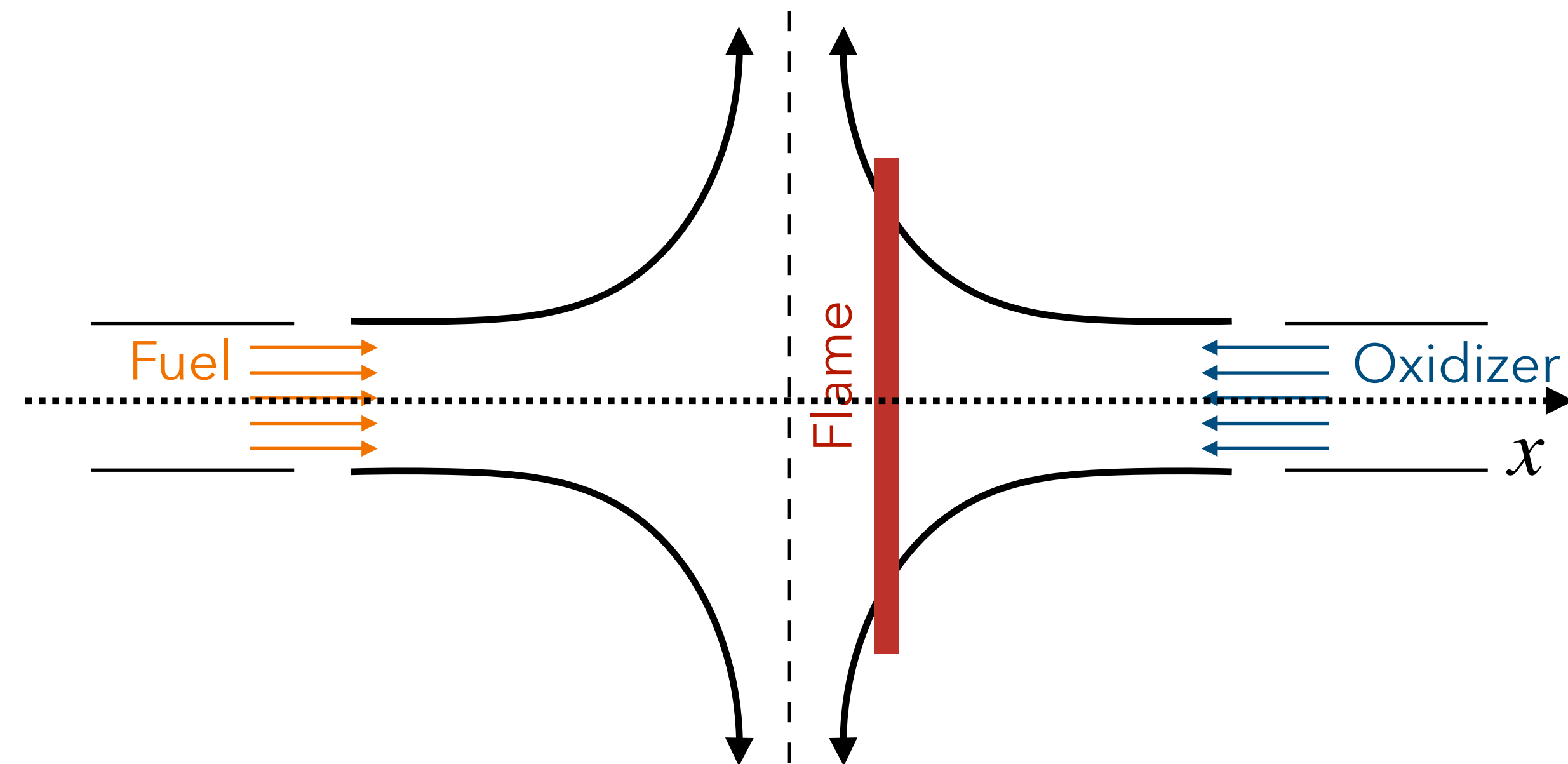
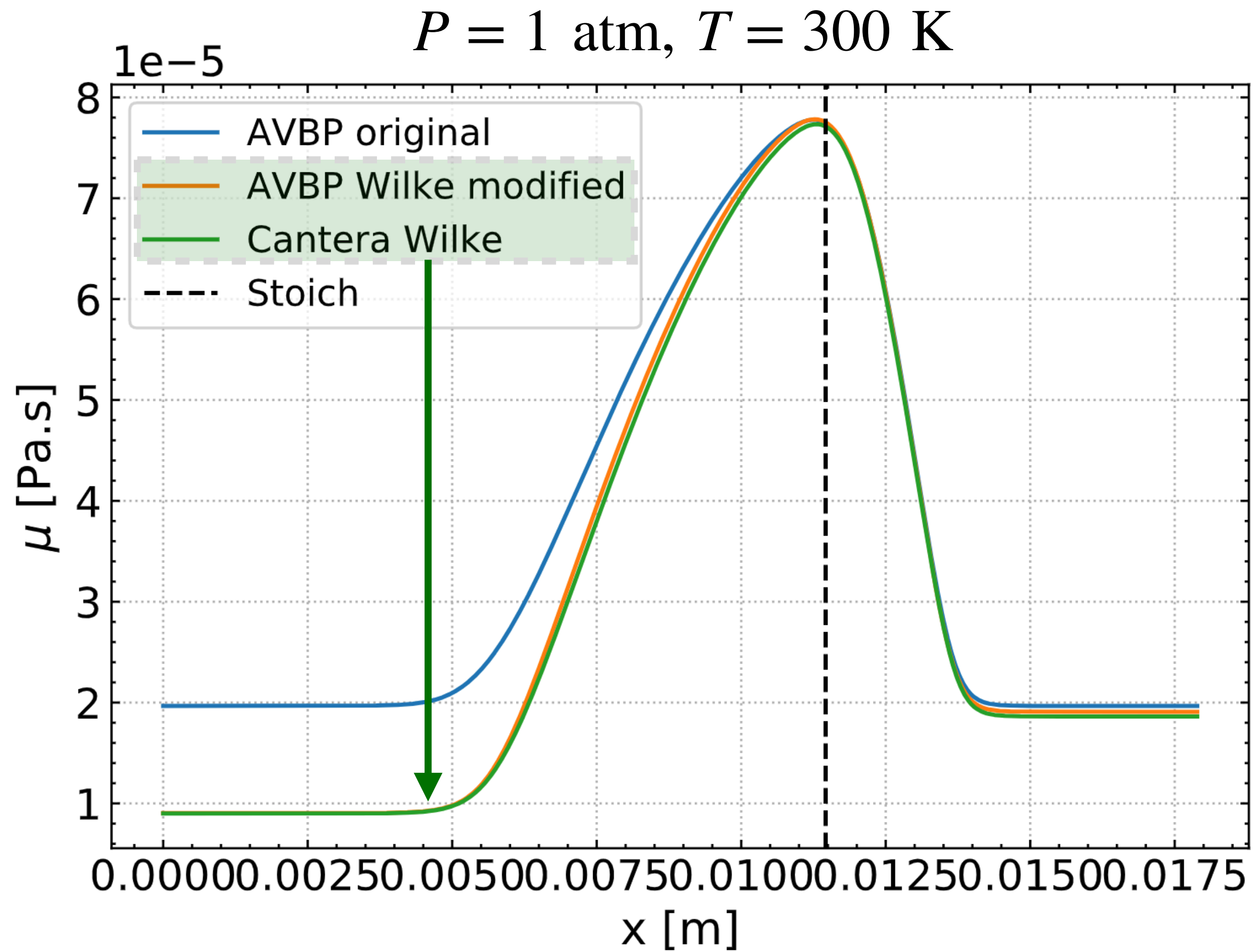
$$\mu_{H_2}(T) = \mu_{ref,H_2} \left( \frac{T}{T_{ref,H_2}} \right)^{\alpha_{H_2}}$$

$$\mu_{air}(T) = \mu_{ref,air} \left( \frac{T}{T_{ref,air}} \right)^{\alpha_{air}}$$



# Transport properties

## Laminar viscosity of the gas

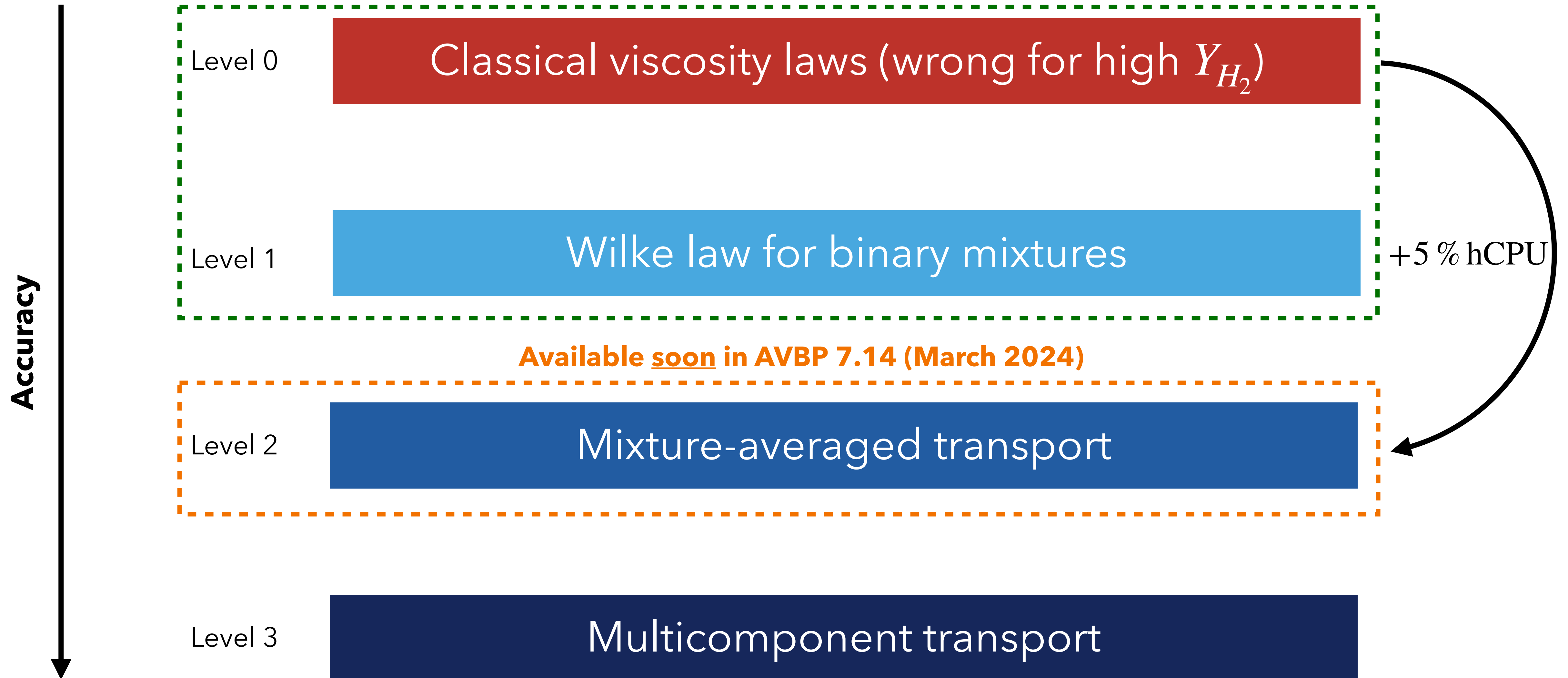




# Transport properties

## Laminar viscosity

Available in AVBP





# Transport properties

## Species diffusion & heat conduction

### Species diffusion coefficient

- Based on Schmidt number

$$D_k = \frac{\mu}{\rho Sc_k}$$

- For Wilke's law: we changed  $\mu$ ,  $Sc_k$  has to be adjusted to recover the correct Species diffusion coefficient

### Heat conduction coefficient

- Based on Prandtl number

$$\lambda = \frac{\mu C_p}{Pr}$$

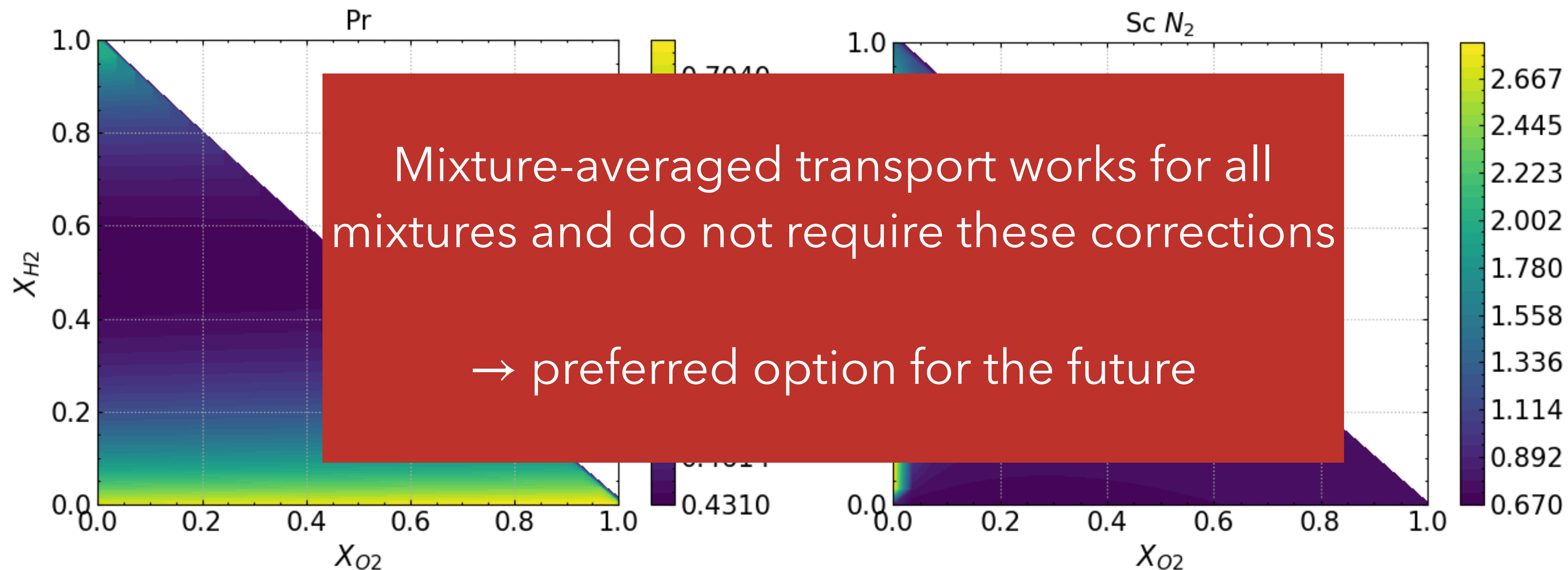
- For Wilke's law: we changed  $\mu$ ,  $Pr$  has to be adjusted to recover the correct heat condition coefficient



# Transport properties

## Wilke's Law: variable Prandlt & Schmidt numbers

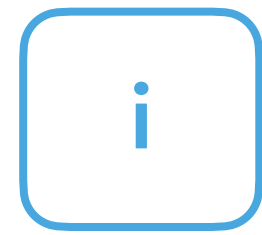
- $Pr_k$  and  $Sc_k$  of the major species are tabulated:  $H_2$ ,  $O_2$ ,  $N_2$ ,  $H_2O$





# Non-adiabatic wall treatments

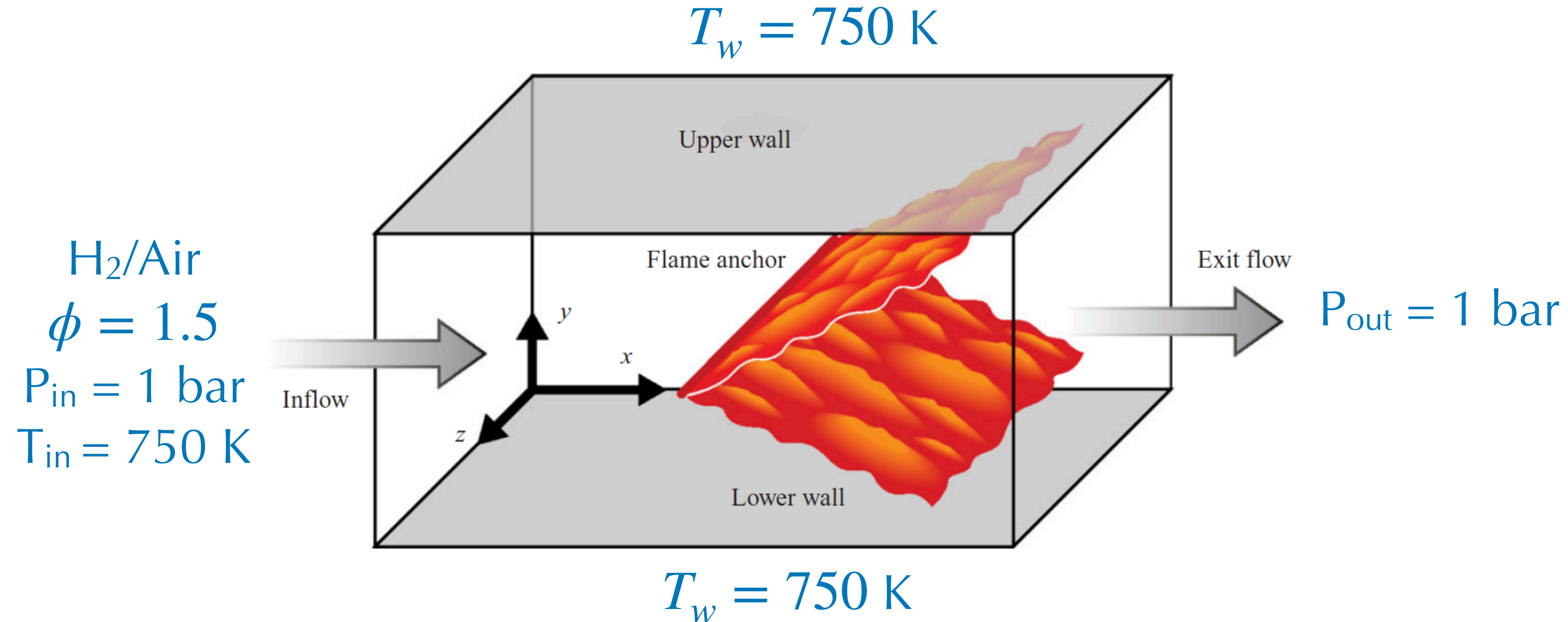
# IFHC model for H<sub>2</sub> FWI



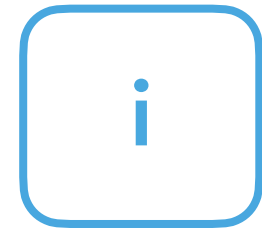
Laminar FWI - standard wall treatment



► A well-known case in the literature of H<sub>2</sub> FWI [1]



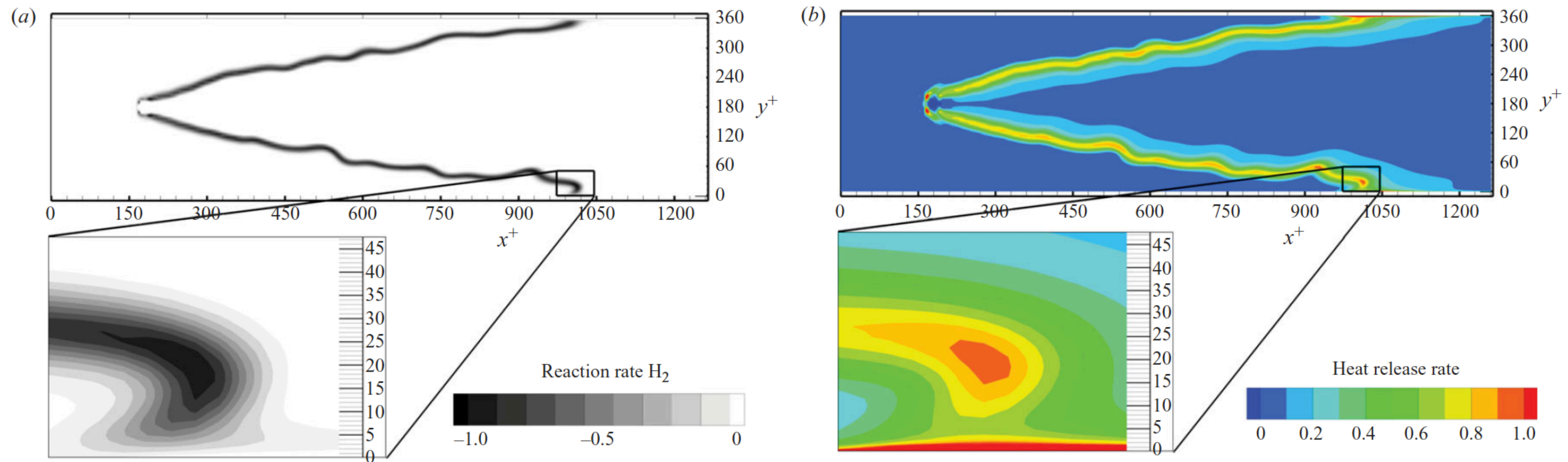
# IFHC model for H<sub>2</sub> FWI



Laminar FWI - standard wall treatment



► A well-known case in the literature of H<sub>2</sub> FWI [1]



a) H<sub>2</sub> reaction rate vanishes at the wall → local quenching

b) Heat release rate peaked values at the wall → unexpected behavior

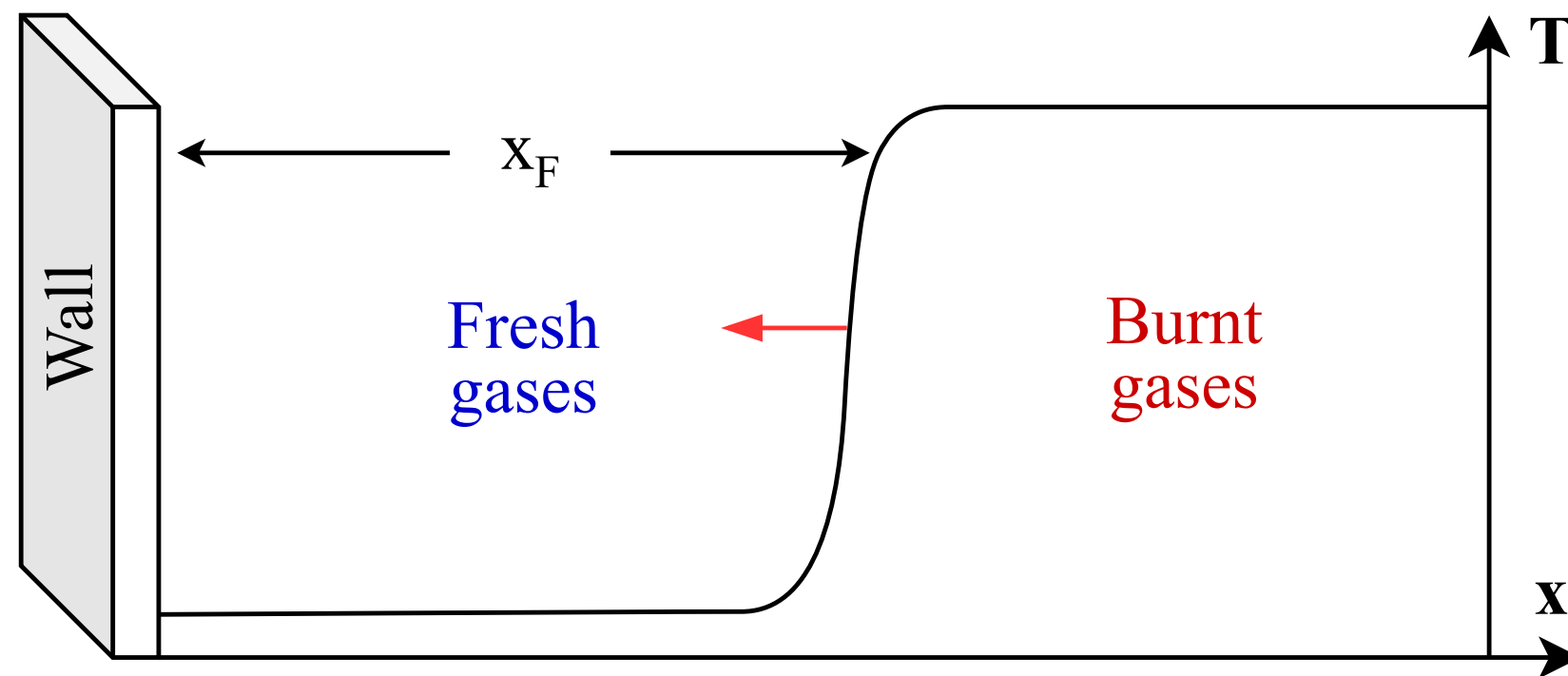
# IFHC model for H<sub>2</sub> FWI

i Laminar FWI - standard wall treatment

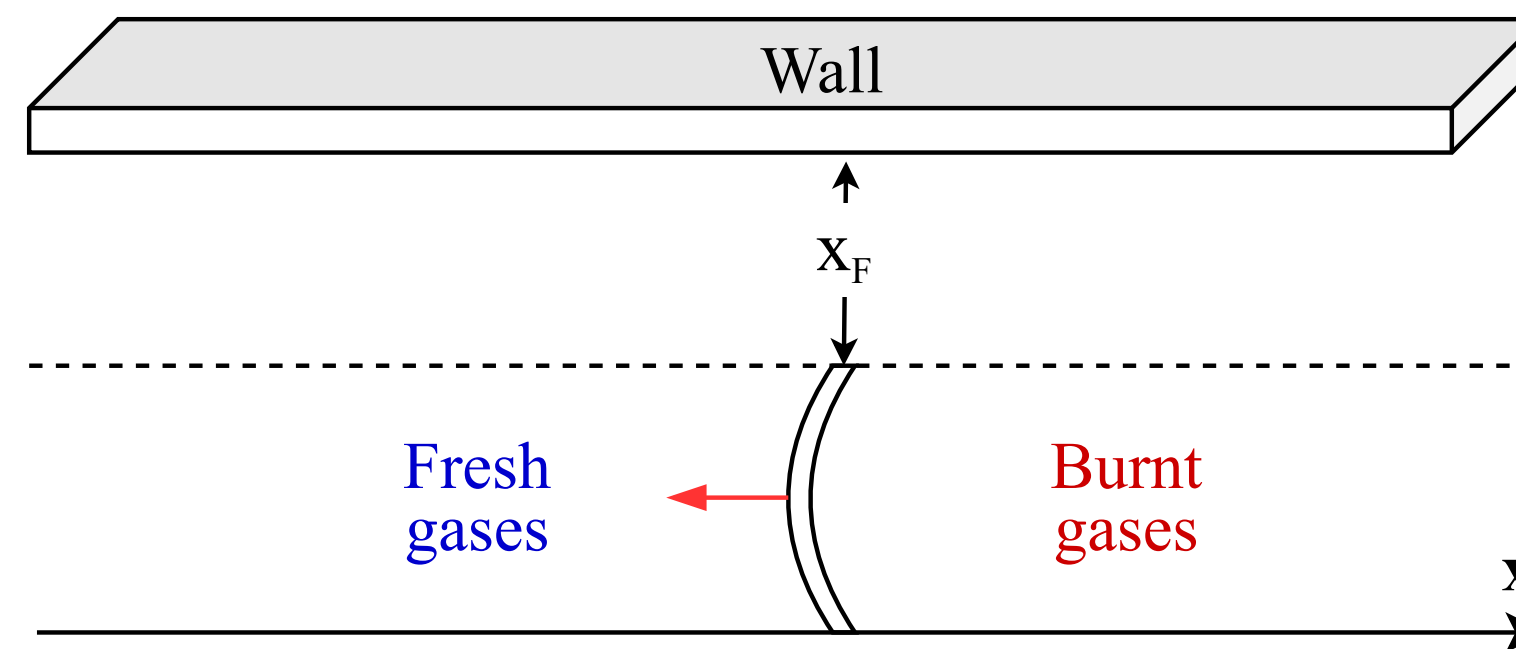


► FWI: complex and transient → worthwhile considering canonical cases [2]

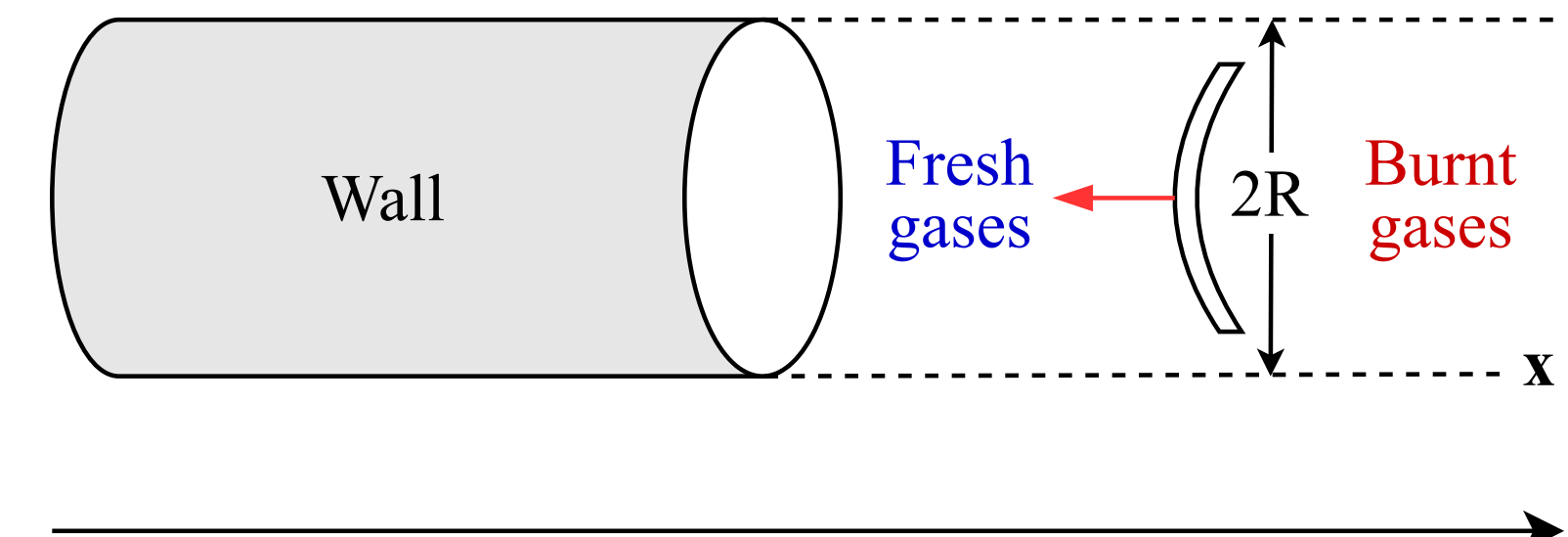
a) Head-On Quenching (HOQ)



b) Side-Wall Quenching (SWQ)



c) Tube quenching

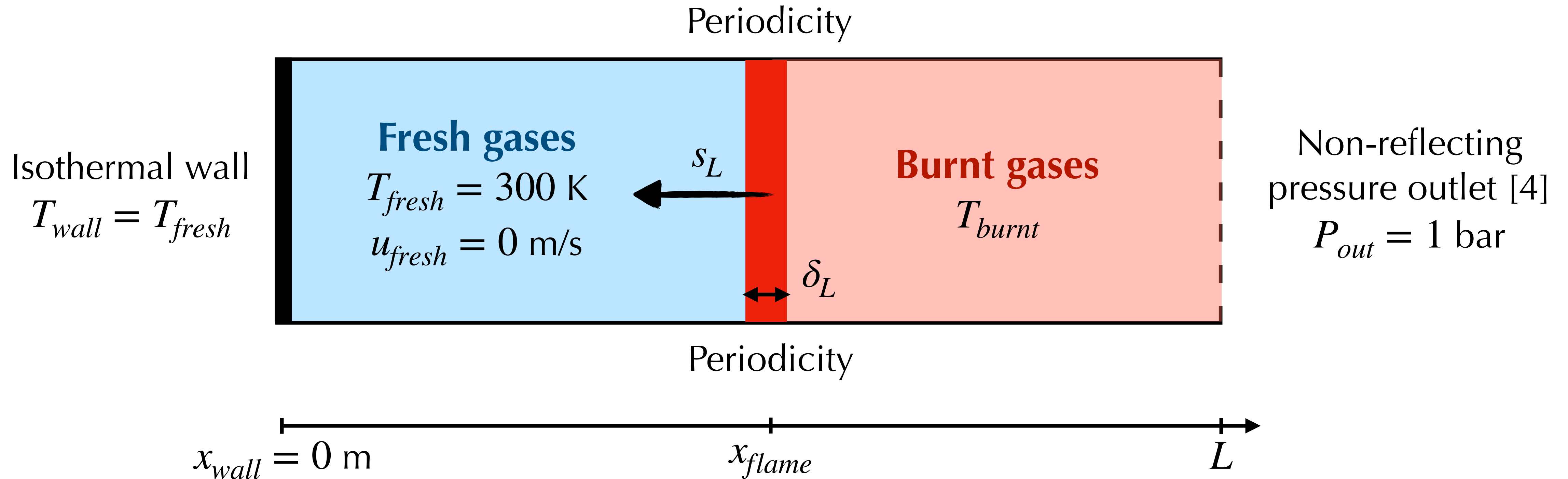




# IFHC model for H<sub>2</sub> FWI

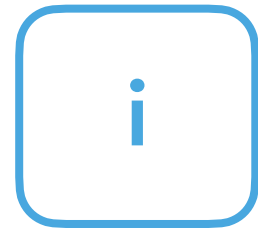
i Laminar FWI - standard wall treatment

► 1D HOQ simulation with **AVBP** [3]



Performed with both **San Diego** [5] and **Burke** [6] schemes

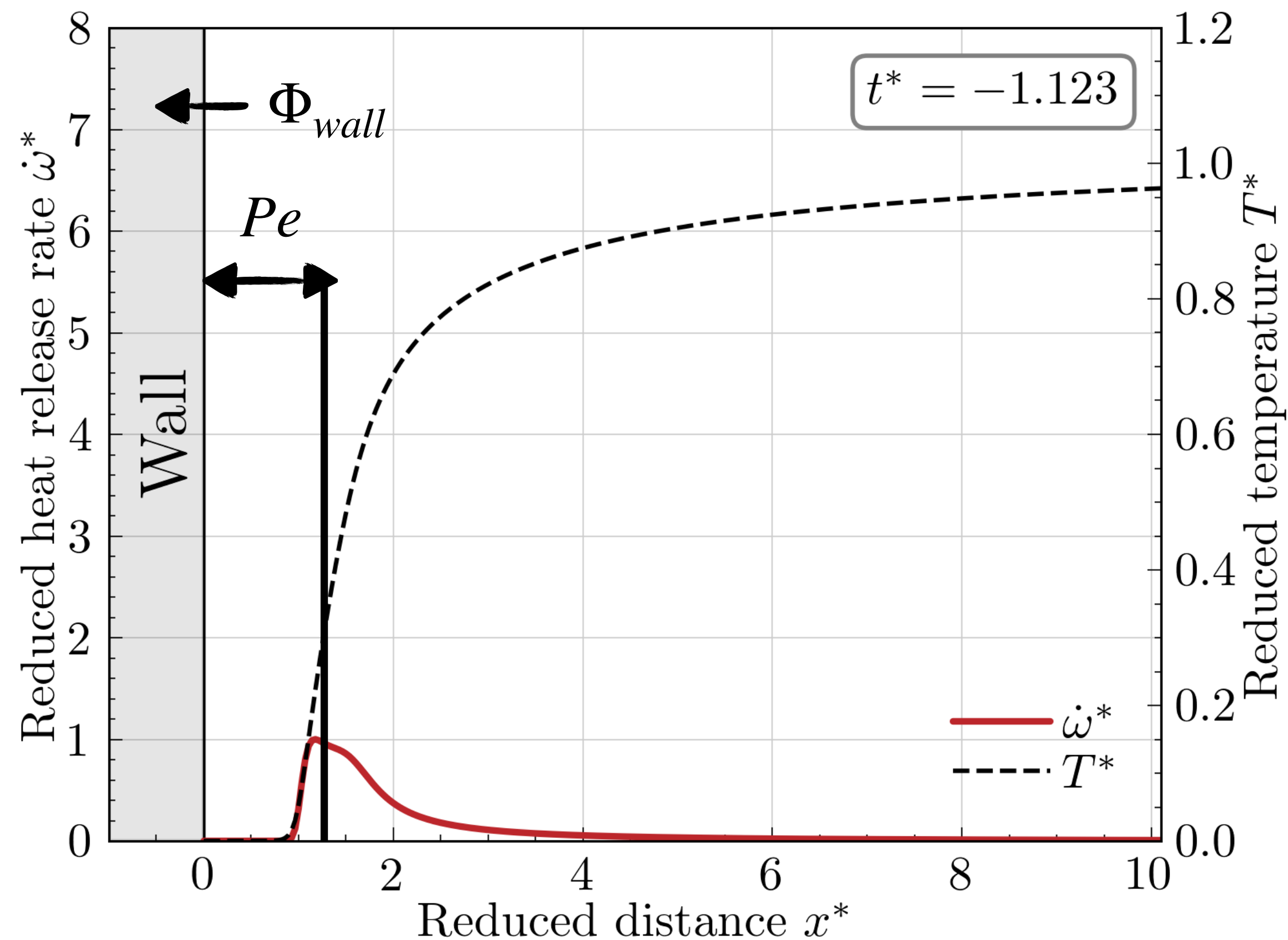
# IFHC model for H<sub>2</sub> FWI



Laminar FWI - standard wall treatment



► Definitions



## ► Reduced variables

- $x^* = x/\delta_L \longrightarrow x^* = 0$ : wall position
- $t^* = (t - t_Q)s_L/\delta_L \longrightarrow t^* = 0$ : quenching instant
- $T^* = (T - T_f)/(T_b - T_f) \longrightarrow T^* = 0$ : fresh gases
- $\dot{\omega}^* = \dot{\omega}_T/\dot{\omega}_T^0 \longrightarrow \dot{\omega}^* = 1$ : free propagation
- $T^* = 1$ : burnt gases
- $t^* < 0$ : free propagation
- $t^* > 0$ : post interaction

## ► Quantities of interest

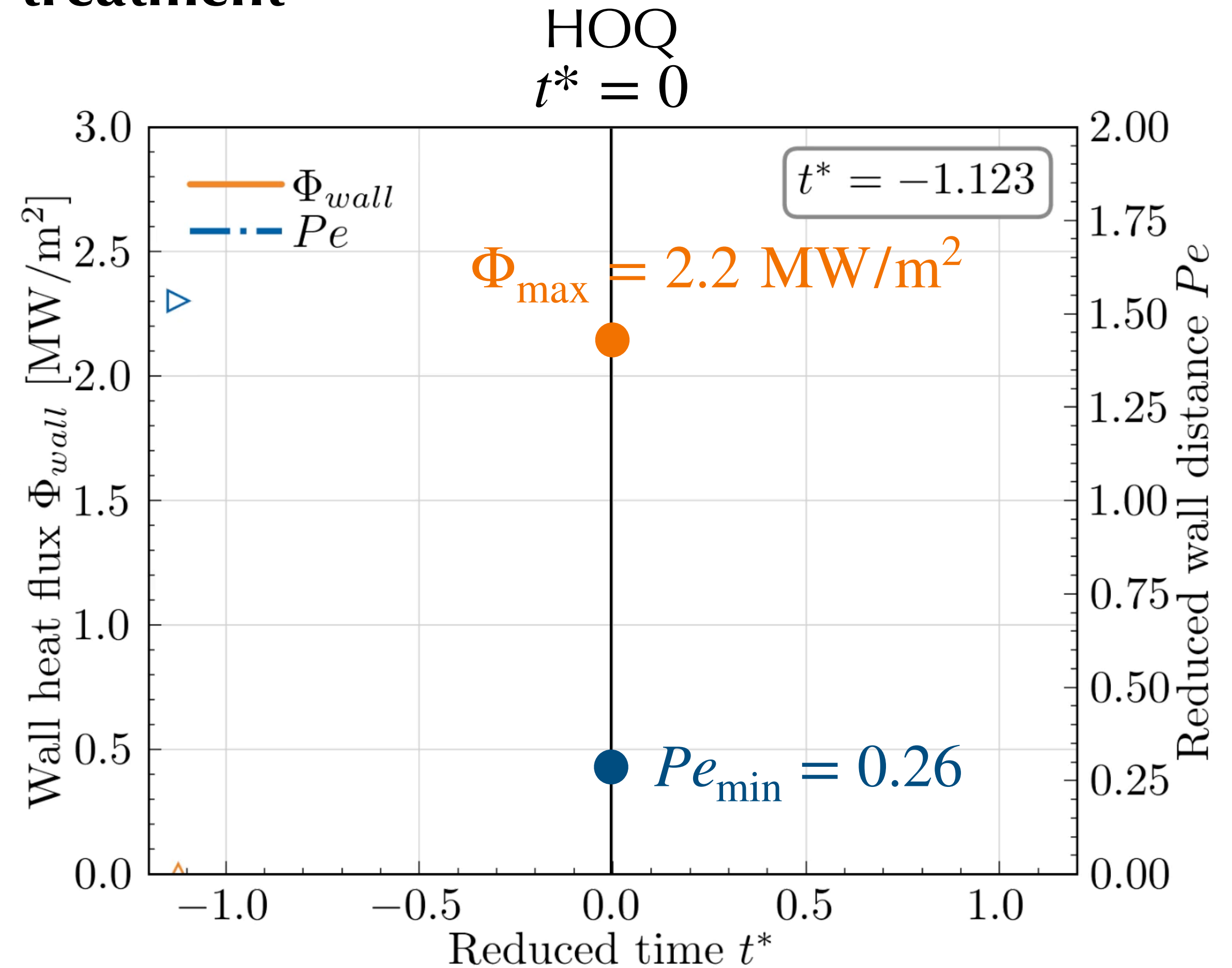
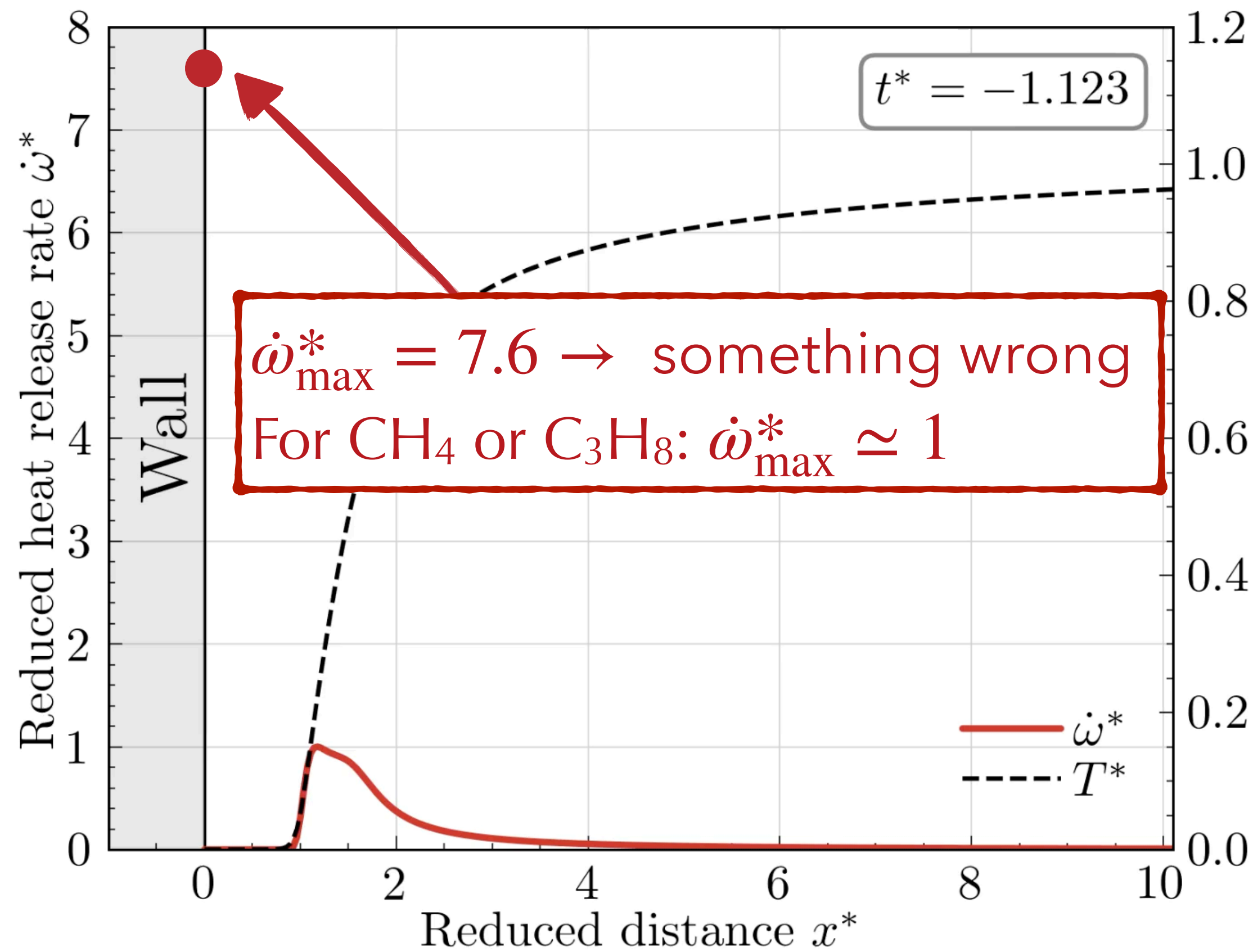
$$Pe = x_{flame}/\delta_L$$

$\Phi_{wall}$   $\longrightarrow$  easiest-to-access physical parameter in EXP.

# IFHC model for H<sub>2</sub> FWI

i Laminar FWI - standard wall treatment

🔥 ▶ Results for a standard inert wall treatment



[1] Gruber et al. (JFM, 2010)

[7] Owston et al. (IJHE, 2007)

[9] Zhao et al. (CNF, 2022)

[11] Lai et al. (IJHFF, 2022)

[8] Mari et al. (CNF, 2016)

[10] Dabireau et al. (CNF, 2003)

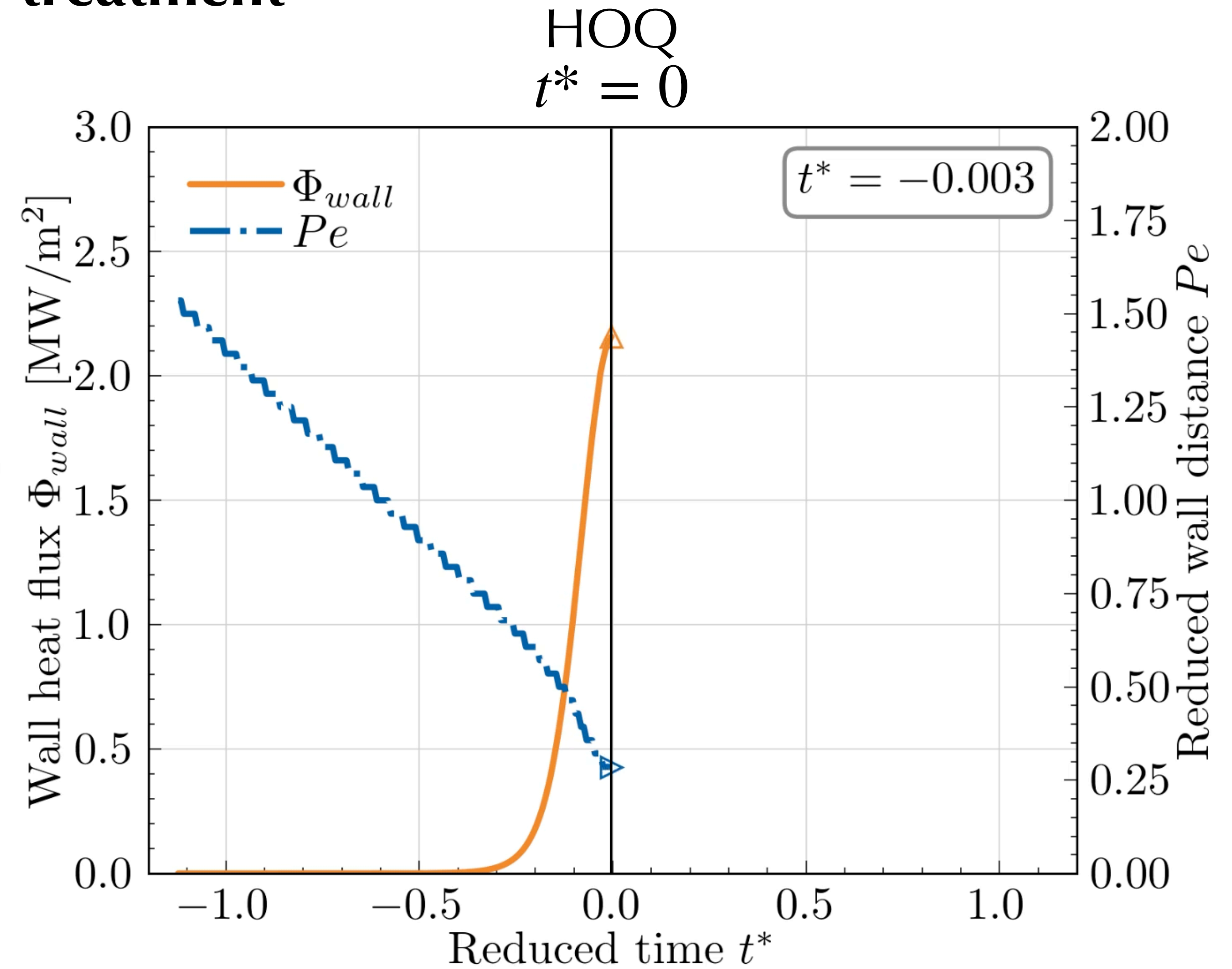
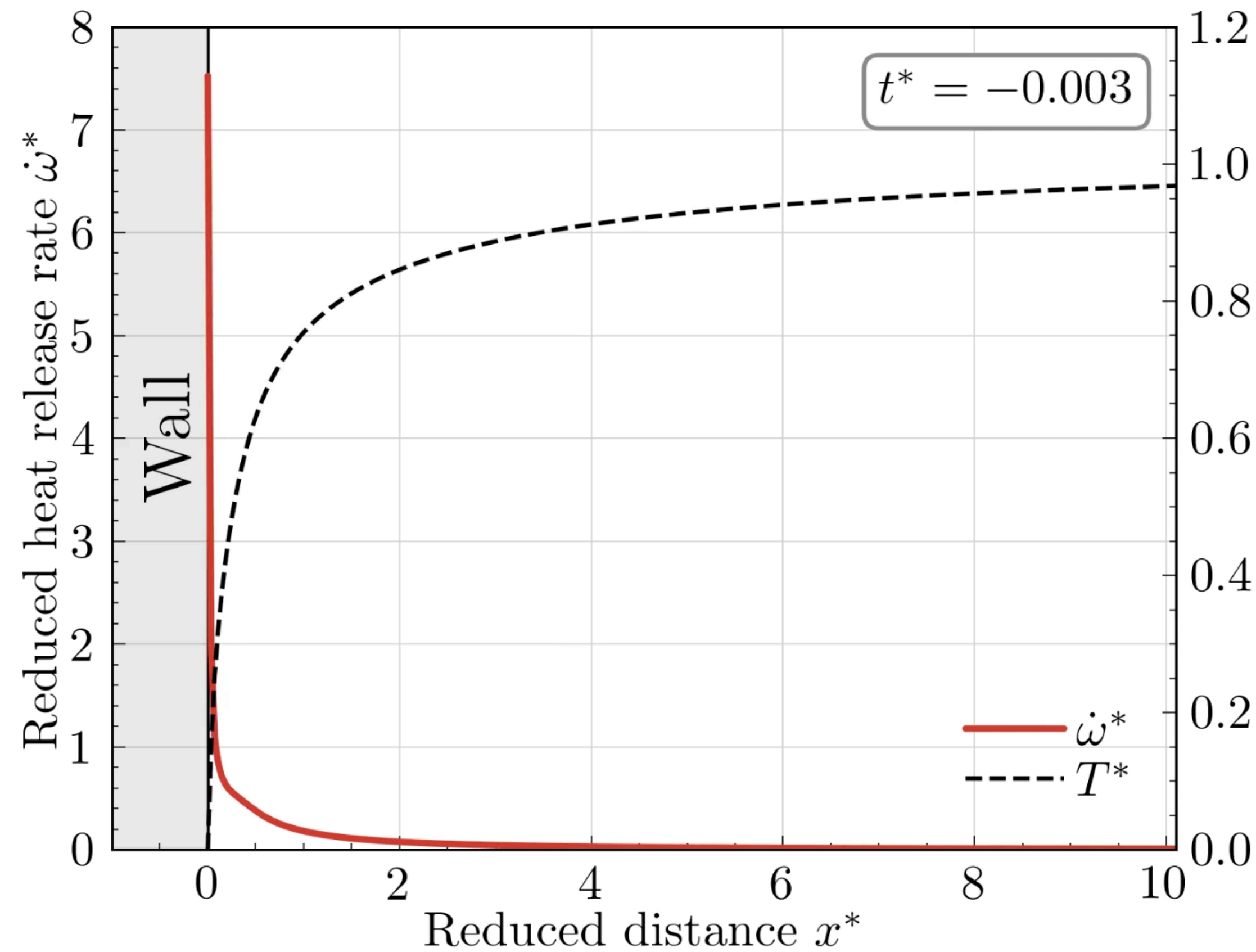
[12] Popp and Baum (CNF, 1997)

[21] De Nardi et al. (CNF, 2024)

# IFHC model for H<sub>2</sub> FWI

i Laminar FWI - standard wall treatment

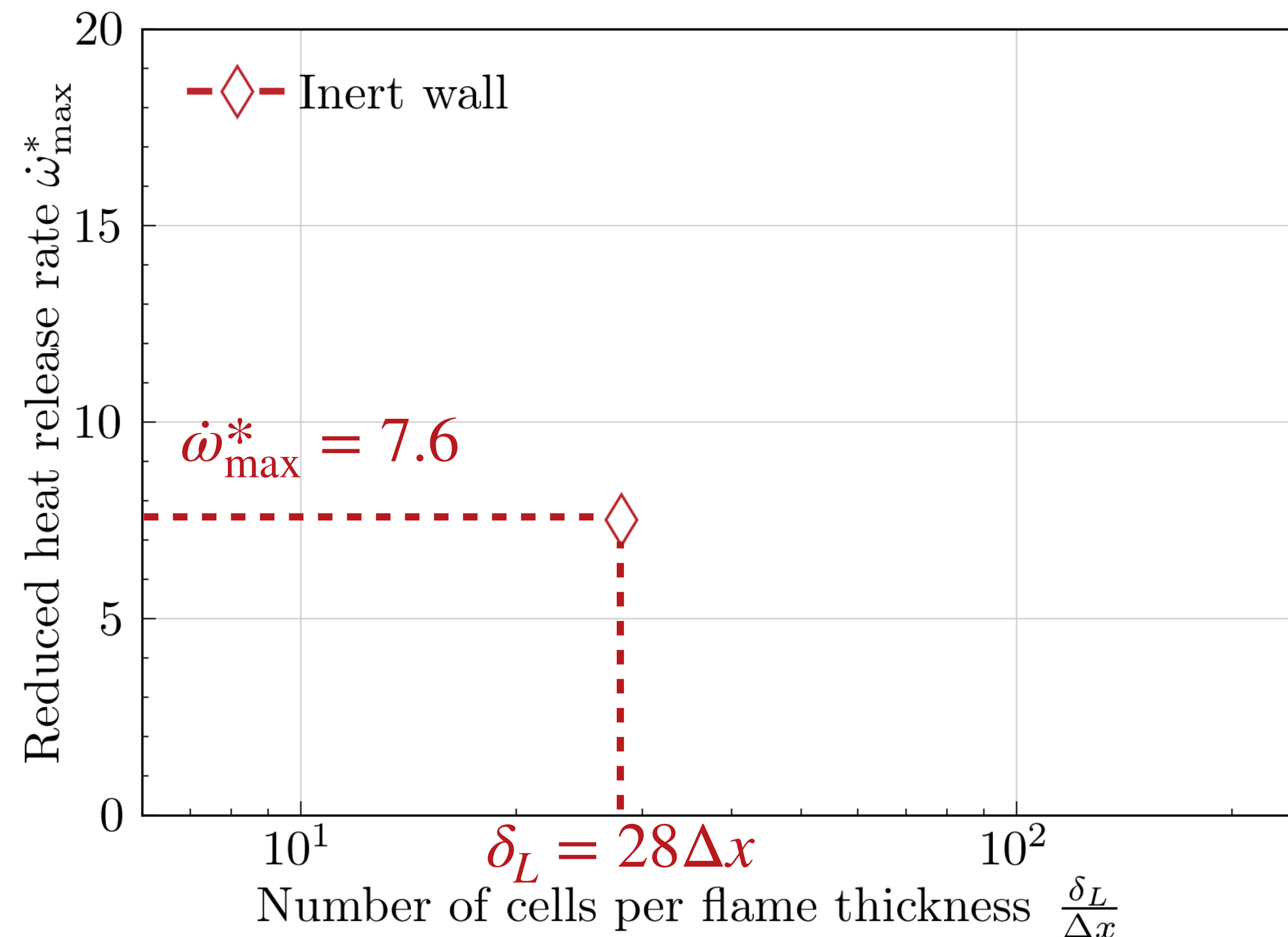
► Results for a standard inert wall treatment



# IFHC model for H<sub>2</sub> FWI

i Laminar FWI - standard wall treatment

► Results for a standard inert wall treatment



- Gradient stiffened during FWI:  
 $\delta_L = 28\Delta x$  recommended [13,15]
- Effect of grid resolution:  
 $\delta_L = 7\Delta x - 224\Delta x$
- $\dot{\omega}^*$  peak at quenching does not converge when the mesh is refined

→ **Ill-posed problem**

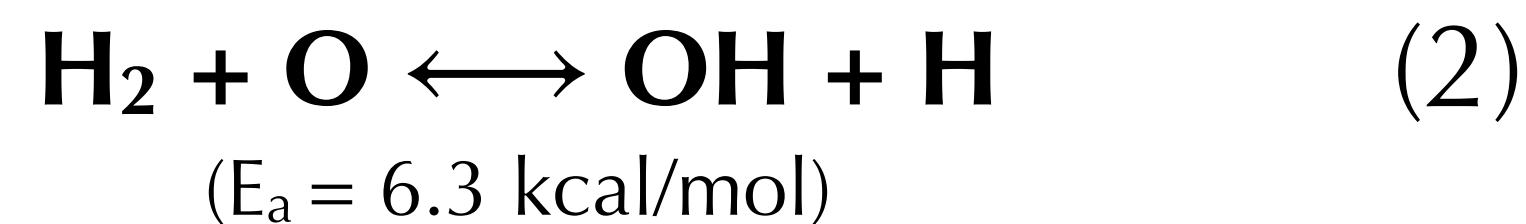
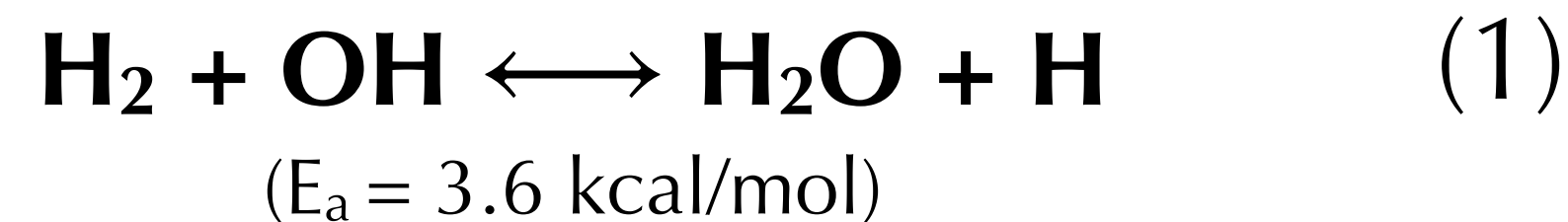
# IFHC model for H<sub>2</sub> FWI

i Laminar FWI - standard wall treatment

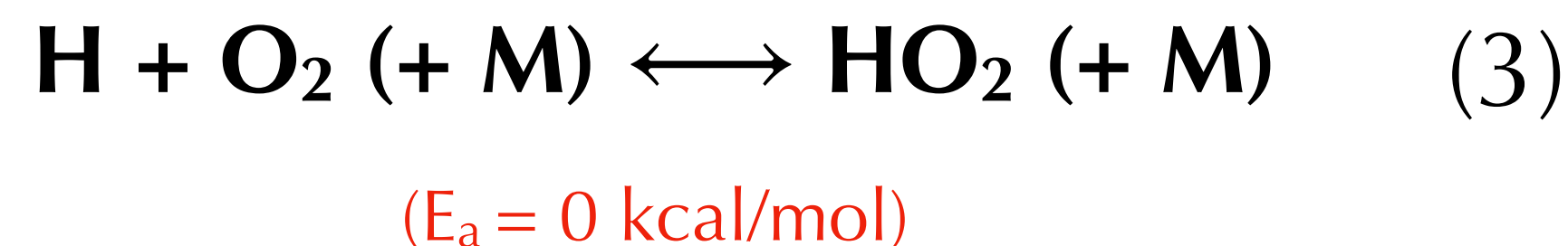
 ► Results for a standard inert wall treatment

## Far from the non-adiabatic wall

- Chain branching reactions:



- Recombination reactions:



## Close to the non-adiabatic wall

- Cooling of the preheat zone → rate of (1) and (2) decreases
- Very fast diffusion of H radical** seems to be the problem ( $Le_H \simeq 0.1$ )
- Only way to consume these radicals is through (3) → **highly exothermic**

# IFHC model for H<sub>2</sub> FWI

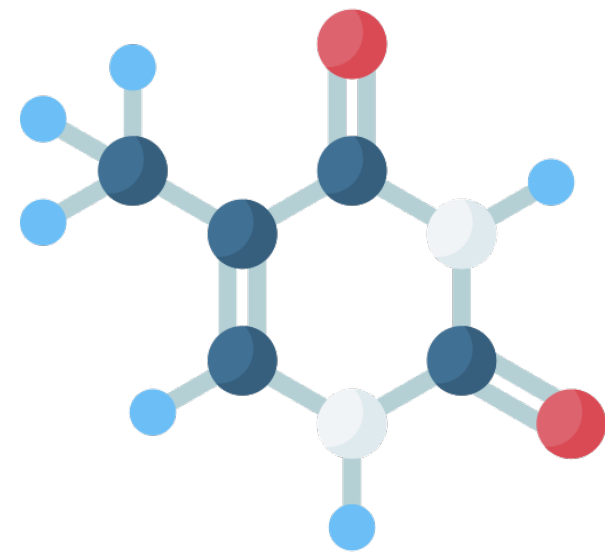
ii

Can we assume wall as inert?



## ► Inert wall assumption?

- Lack of experimental validation for premixed H<sub>2</sub> FWI ( $\Phi_{wall}$ )
- Sticking coefficient for H/O/OH to match EXP. explosion limits [20, 21]
- Wall-coating material type → significant impact on FWI [22]



## ► Heterogeneous catalysis: using detailed surface chemistry mechanisms

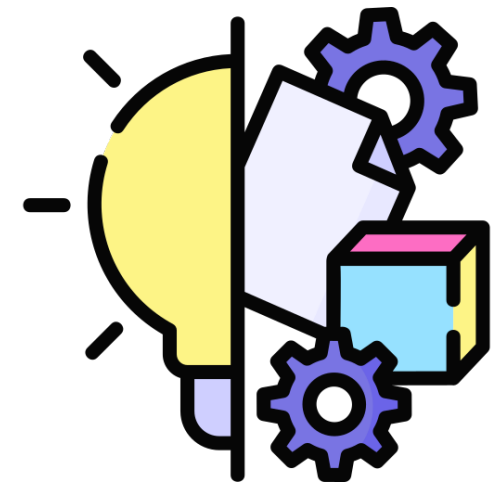
- Exist for hydrocarbons [23] and H<sub>2</sub> [24, 25]
- Also suffer from a lack of experimental validation
- Can significantly increase computational cost!



# IFHC model for H<sub>2</sub> FWI

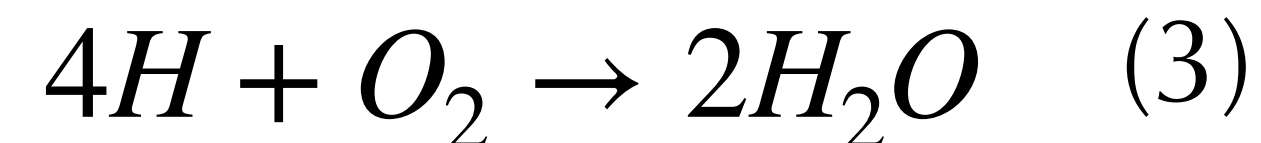
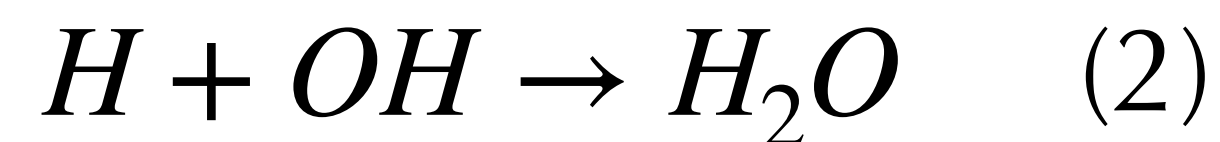
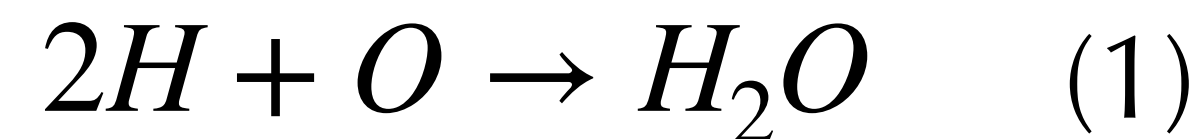
iii

A simplified approach for FWI



## ► Infinitely Fast Heterogeneous Catalysis (IFHC) [21]

- Global surface chemistry: total, irreversible, occurring in 1 timestep:



- Energy conservation:

$$\dot{\omega}_{T,IFHC} = - \sum_k^N \dot{\omega}_{k,IFHC} \times \Delta h_{f,k}^o$$

Heat released at wall  
by IFHC

Surface reaction  
rate of species k

Enthalpy of formation of species k

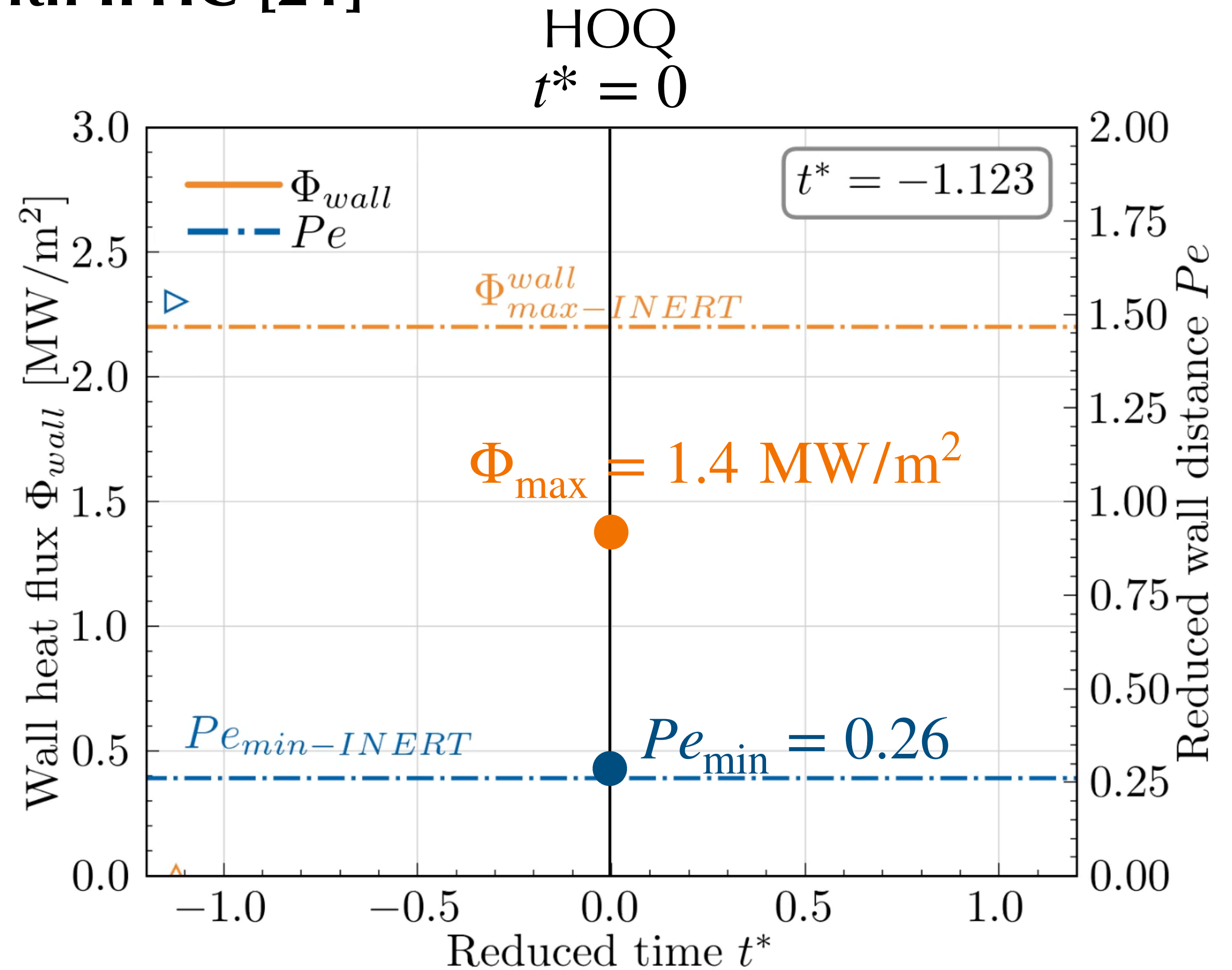
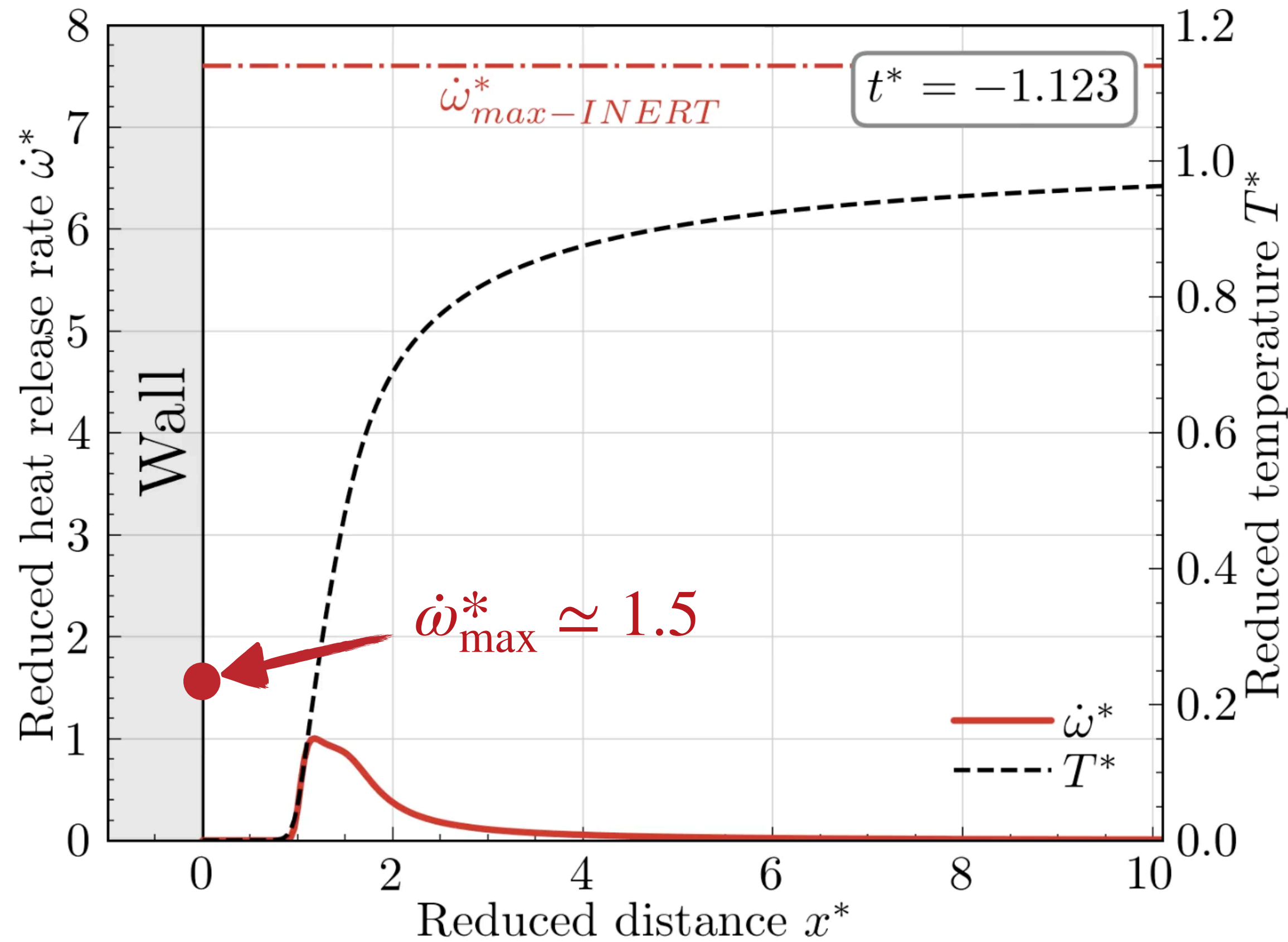


# IFHC model for H<sub>2</sub> FWI

iv Laminar FWI - catalytic wall using IFHC



► Results for an isothermal wall with IFHC [21]

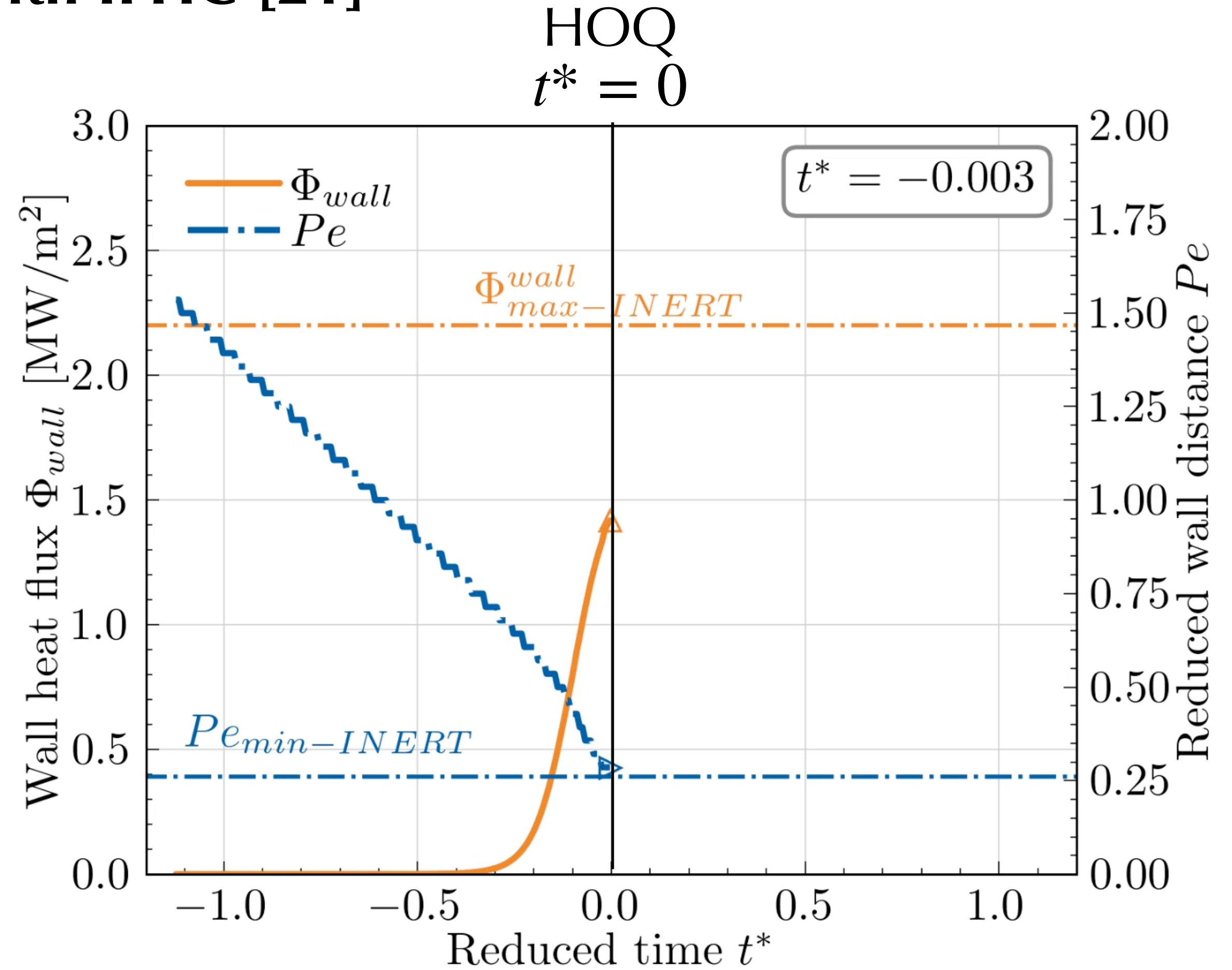
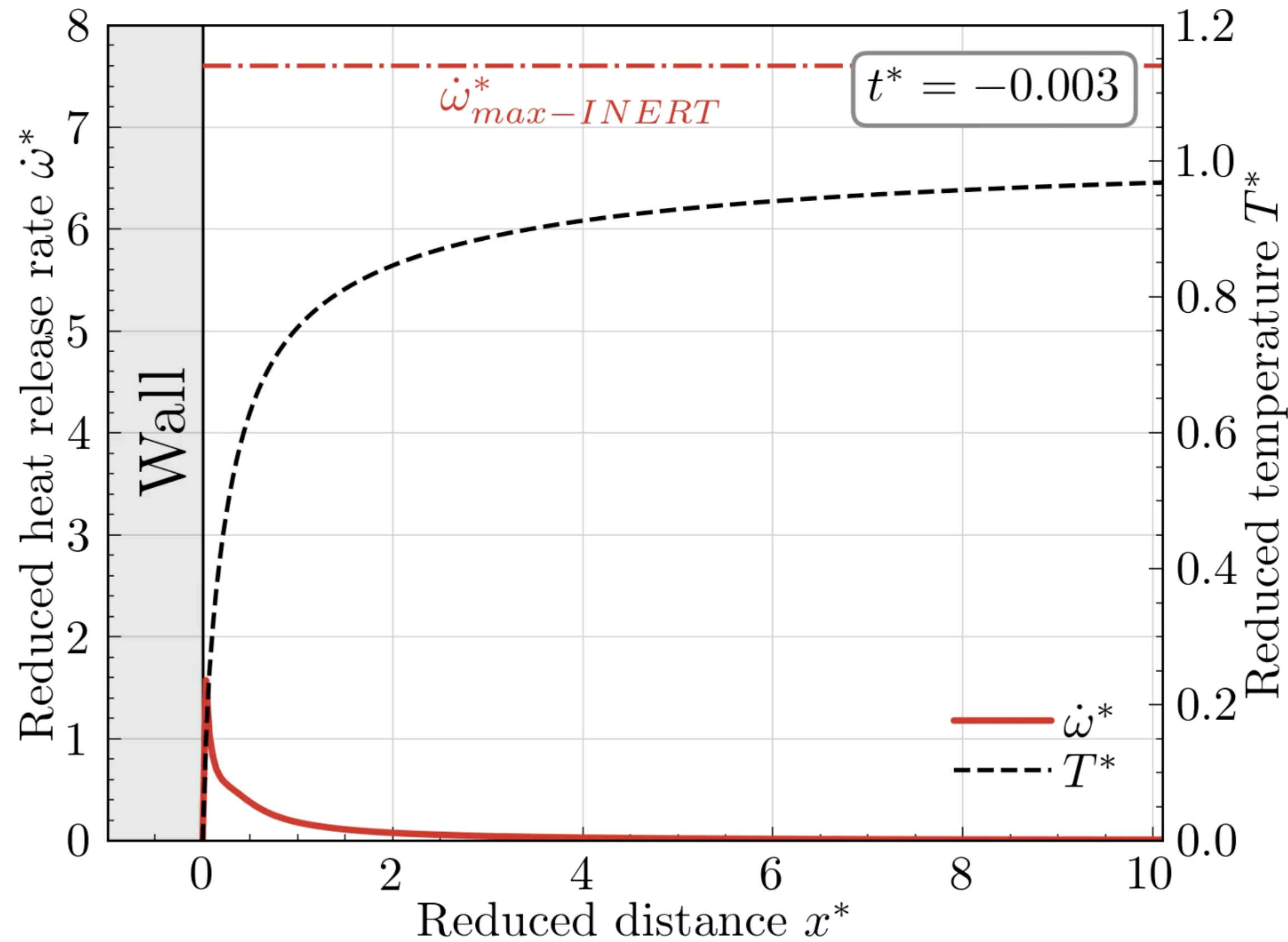


# IFHC model for H<sub>2</sub> FWI

iv Laminar FWI - catalytic wall using IFHC



► Results for an isothermal wall with IFHC [21]



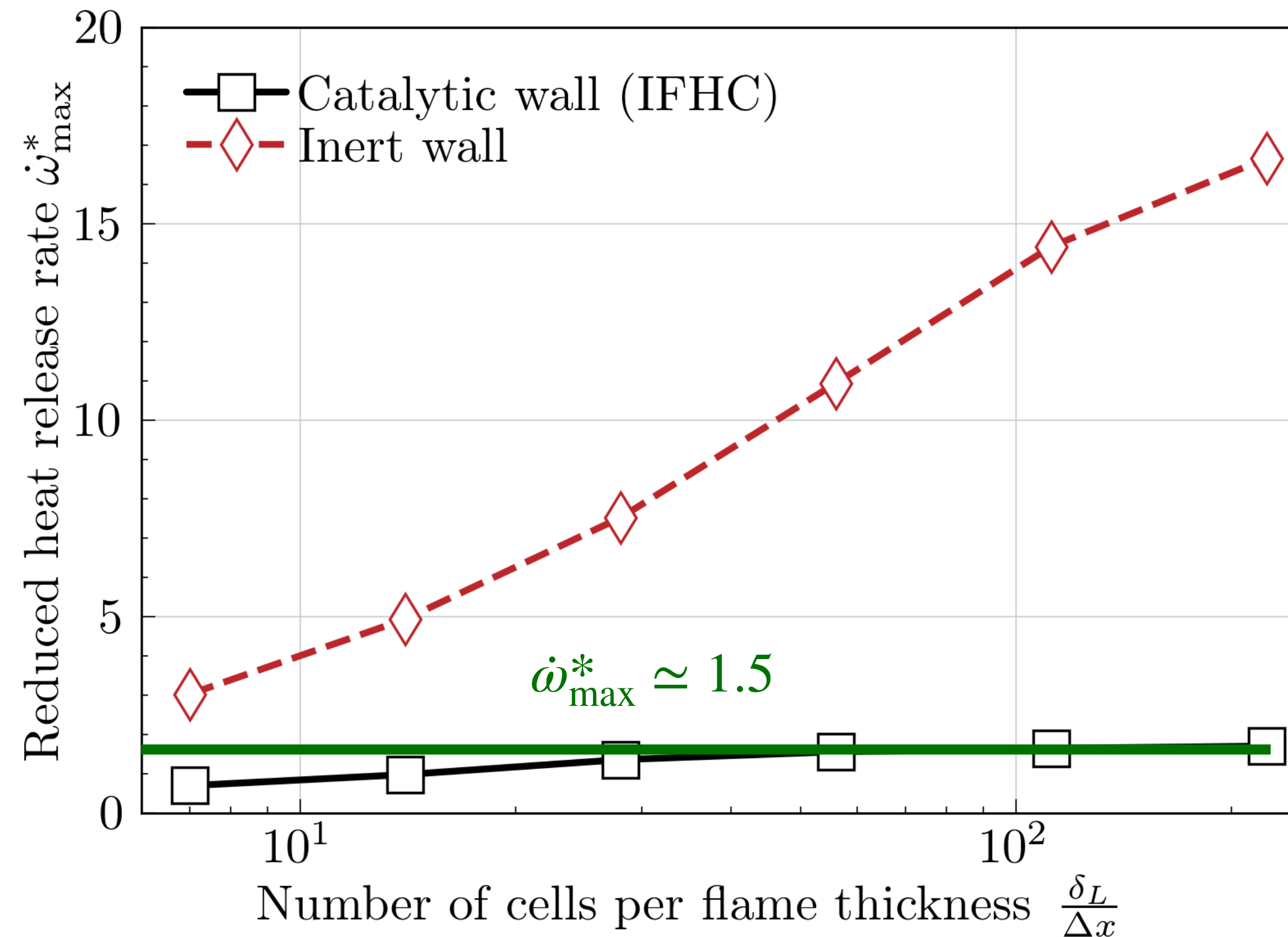
# IFHC model for H<sub>2</sub> FWI

iv

Laminar FWI - catalytic wall using IFHC



► Results for an isothermal wall with IFHC [21]



- Effect of grid resolution with IFHC:  
 $\delta_L = 7\Delta x - 224\Delta x$
- $\dot{\omega}^*$  peak value controlled with IFHC

→ **Grid convergence retrieved**  
at  $\delta_L = 28\Delta x$



# Adaptation of TFLES

# Thickened Flame model: TFLES

## Key parameters

- Balance equation for the mass fraction of species  $k$

$$\frac{\partial \rho Y_k}{\partial t} + \nabla \cdot (\rho \mathbf{u} Y_k) = \nabla \cdot (\rho D_k \nabla Y_k) + \dot{\omega}_k$$

- Balance equation for the mass fraction of species  $k$  in the **TFLES model**

$$\frac{\partial \rho Y_k}{\partial t} + \nabla \cdot (\rho \mathbf{u} Y_k) = \nabla \cdot (\rho D F E \nabla Y_k) + \frac{E}{F} \dot{\omega}_k$$

### $F \equiv$ Thickening factor

- The larger  $F$  is, the less points we need in the flame front
- Thickened flame thermal thickness:

$$\delta_L^1 = F \delta_L^0$$

### $E \equiv$ Efficiency function

- Damkohler number:  $Da^1 = Da^0 / F$
- Karlovitz number:  $Ka^1 = Ka^0 / F$
- Flame front wrinkles less in the presence of turbulence
- $E$  compensate this loss

29

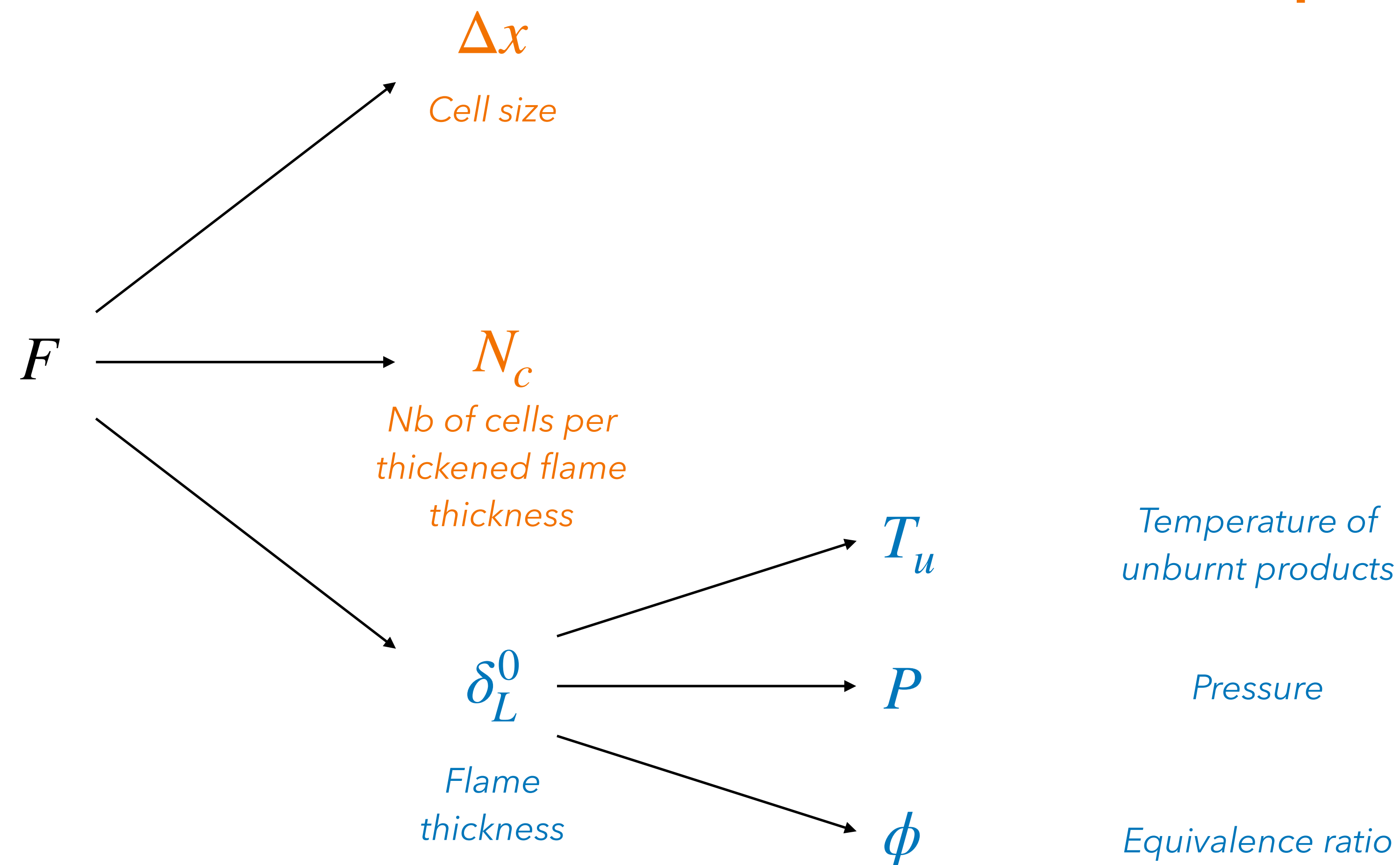


# Thickened Flame model: TFLES

*Key parameters: dependancy of the thickening factor*

**Physical parameters**

**Model parameters**



$$F\delta_L^0 = N_c\Delta x$$

**Dynamic thickening  
(DTFLES)**



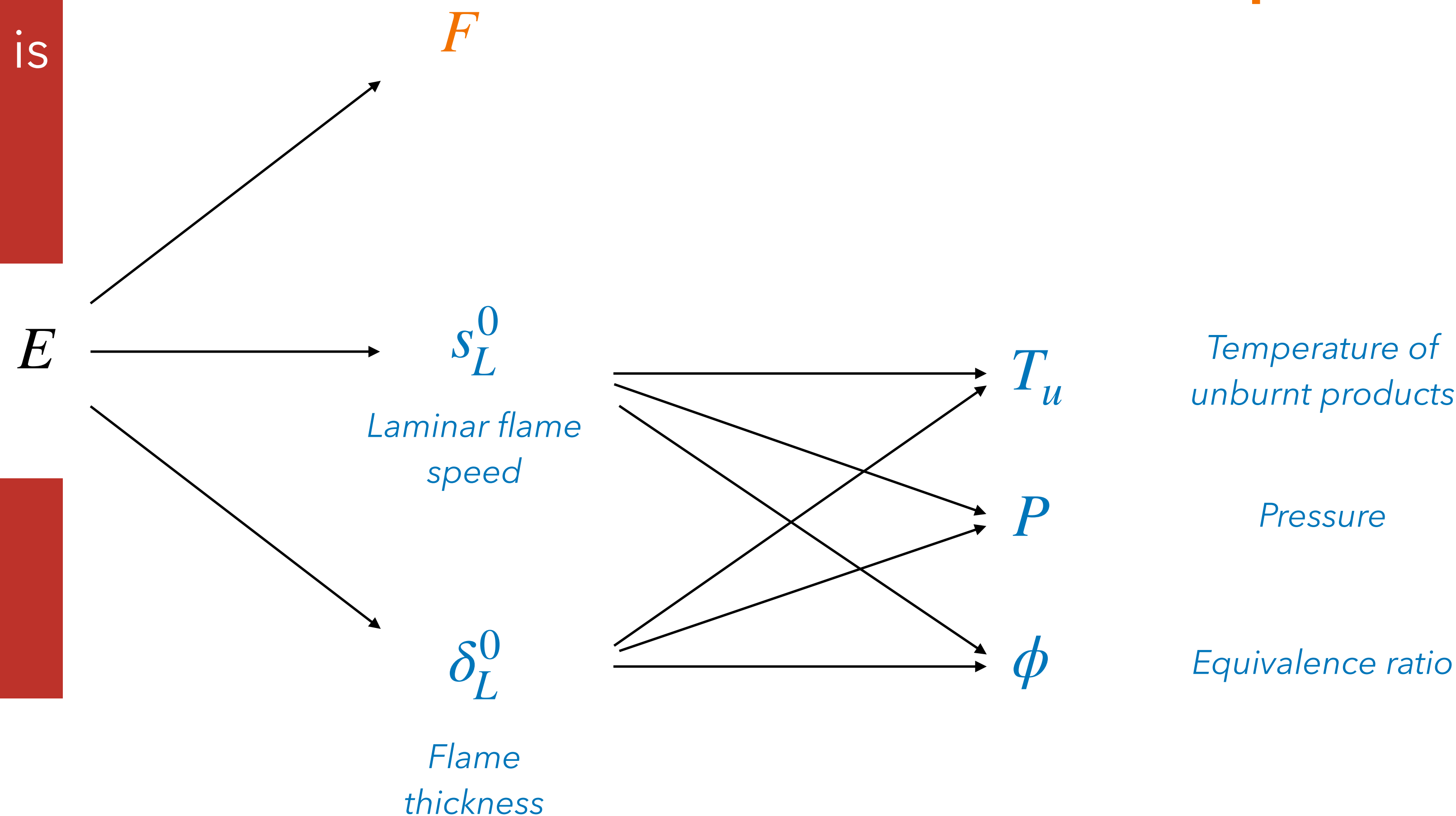
# Thickened Flame model: TFLES

*Key parameters: dependancy of the efficiency function*

Recovering the flame properties in the flame front is **key** to the modeling of thickened flames

We need:  
 $(T_u, P, \phi)$

**Typical model:**  
**Charlette efficiency function**



**Physical parameters**

**Model parameters**

[1] F. Charlette et al. A power-law wrinkling model for LES of premixed turbulent combustion: Part I-non-dynamic formulation and initial tests. *Combustion and Flame*, 131:159–180, 2002.

# Thickened Flame model: TFLES

## Equivalence ratio based on Bilger mixture fraction

Passive scalar constructed from elemental mass fractions ( $Z_p$ ).

$Y_F$  and  $Y_{Ox}$  vary across the flame due to decomposition;  $Z_C$ ,  $Z_H$ , and  $Z_O$  do not.

$$Z_p = \sum_{k=1}^N \frac{a_{p,k} W_p}{W_k} Y_k$$

$W_k$  → Mol. weight of species  $k$   
 $Y_k$  → Mass fraction of species  $k$   
 $W_p$  → Mol. weight of element  $p$   
 $a_{p,k}$  → No. of  $p$  element atoms in species  $k$

$$\beta = \sum_{p=1}^P \gamma_p Z_p$$

Linear combination so that  $\beta_{st} = 0$ .

$$\xi = \frac{\beta - \beta_{Oxid}}{\beta_{Fuel} - \beta_{Oxid}}$$

Normalization between the two bounding cases.

These operations require the definition of the oxidizer composition and the fuel composition.

$$\Phi = \frac{\xi}{(1 - \xi)} * \frac{1 - \xi_{st}}{\xi_{st}}$$

Eq. ratio from mixture fraction.

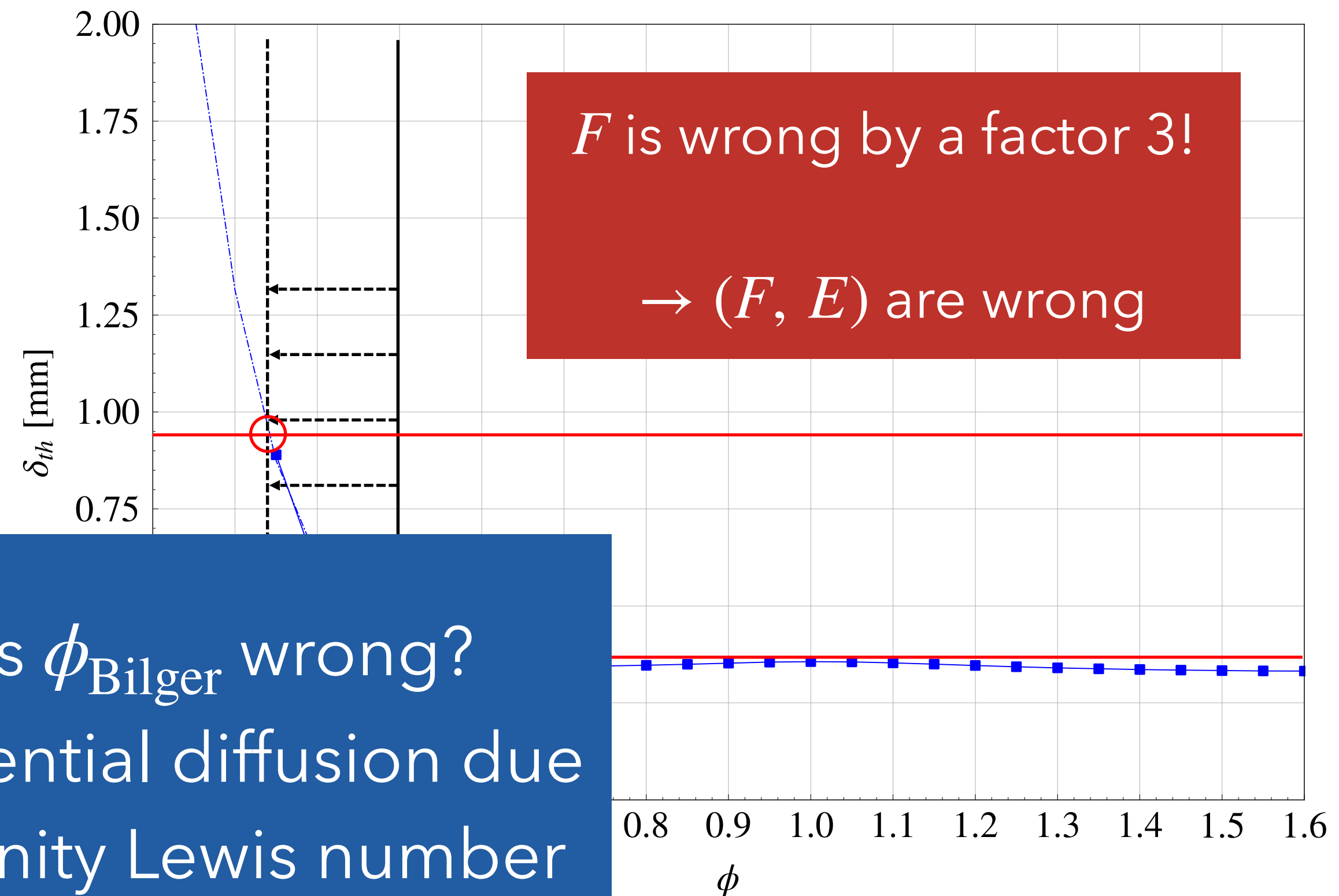
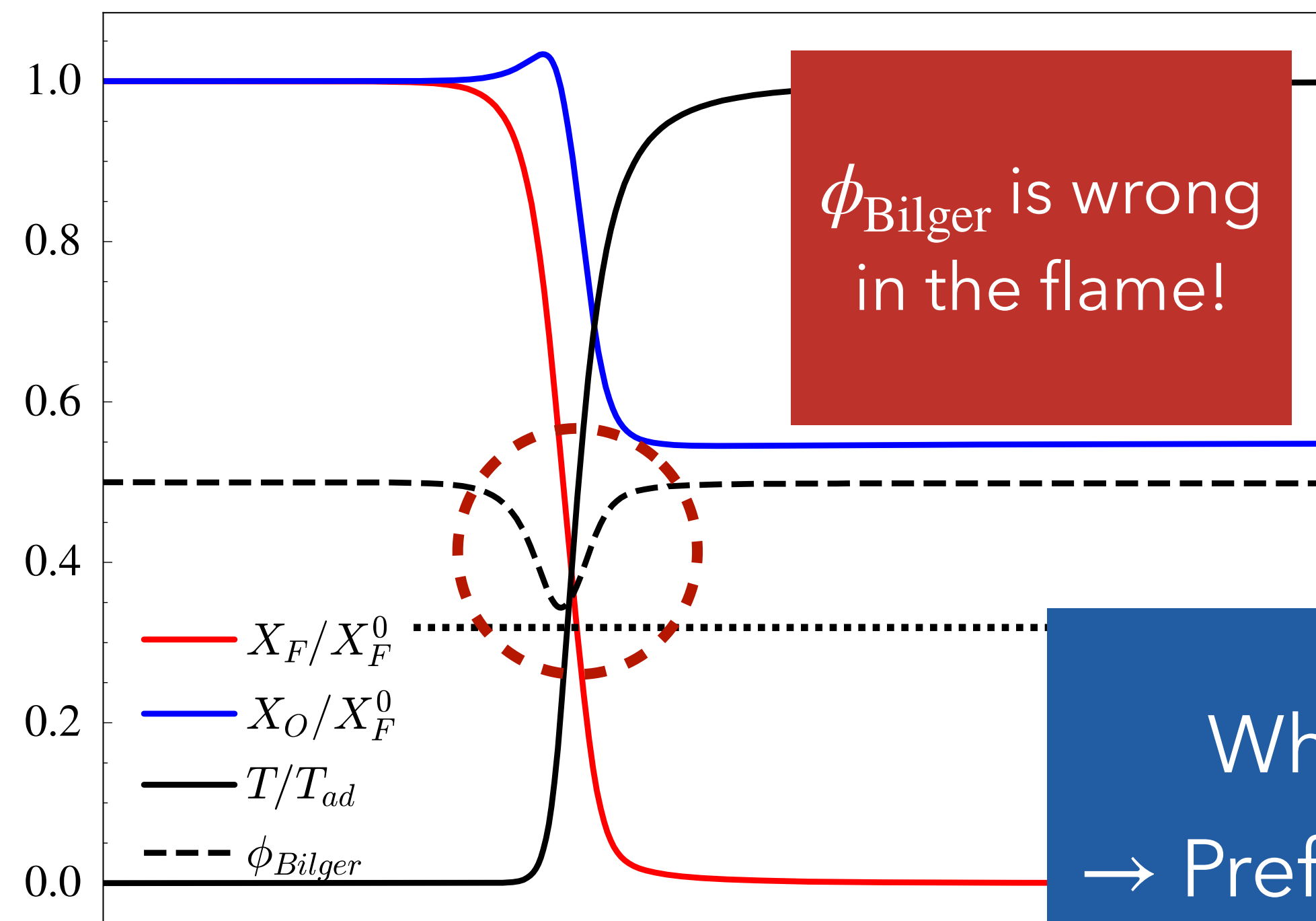
**Classical formulation to retrieve the equivalence ratio**



# Thickened Flame model: TFLES

Equivalence ratio based on Bilger mixture fraction

## Laminar flame thickness



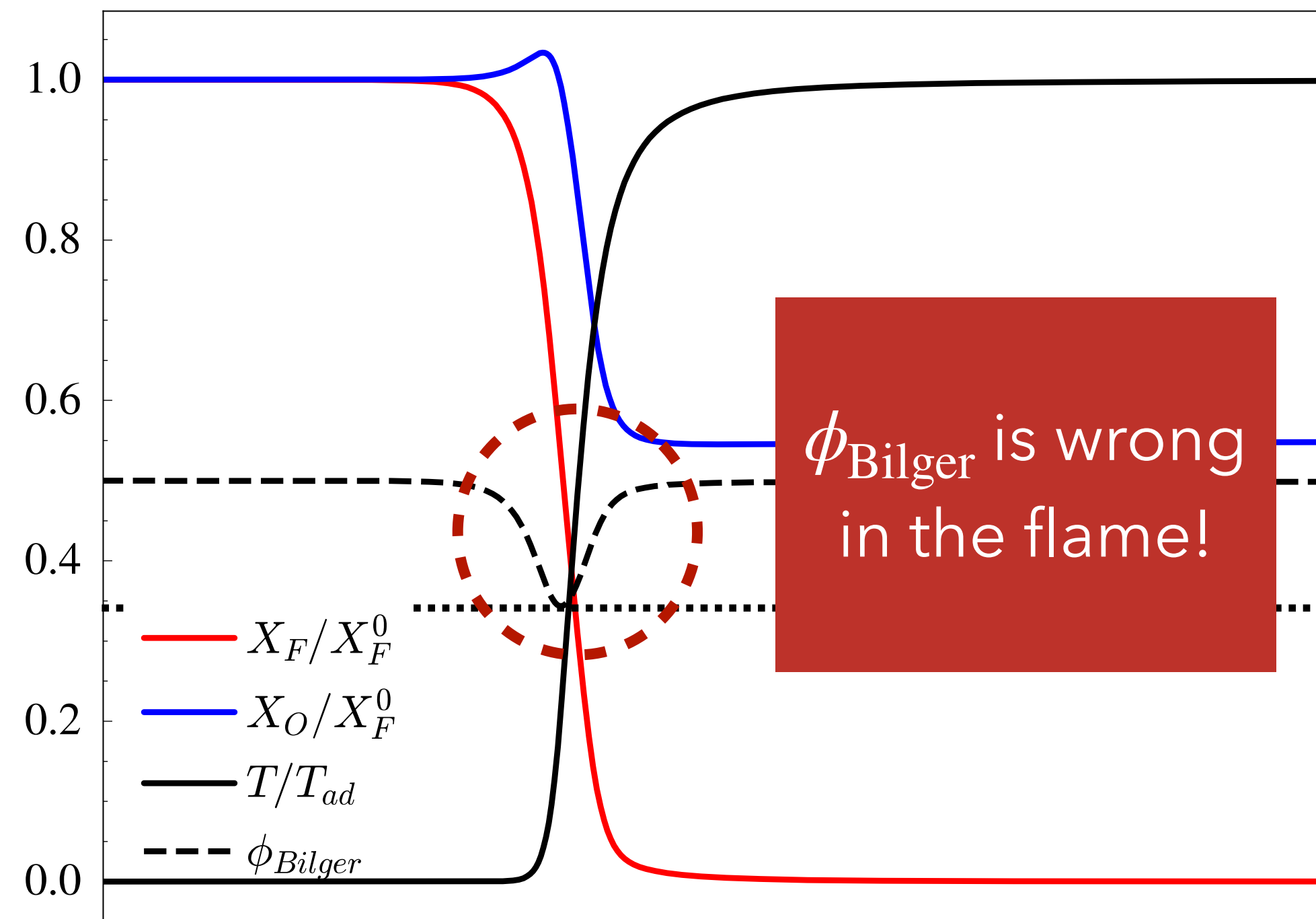
Lean H<sub>2</sub>-Air Flame ( $\phi = 0.5$ ,  $T_u = 300$  K,

Why is  $\phi_{Bilger}$  wrong?  
→ Preferential diffusion due to non-unity Lewis number

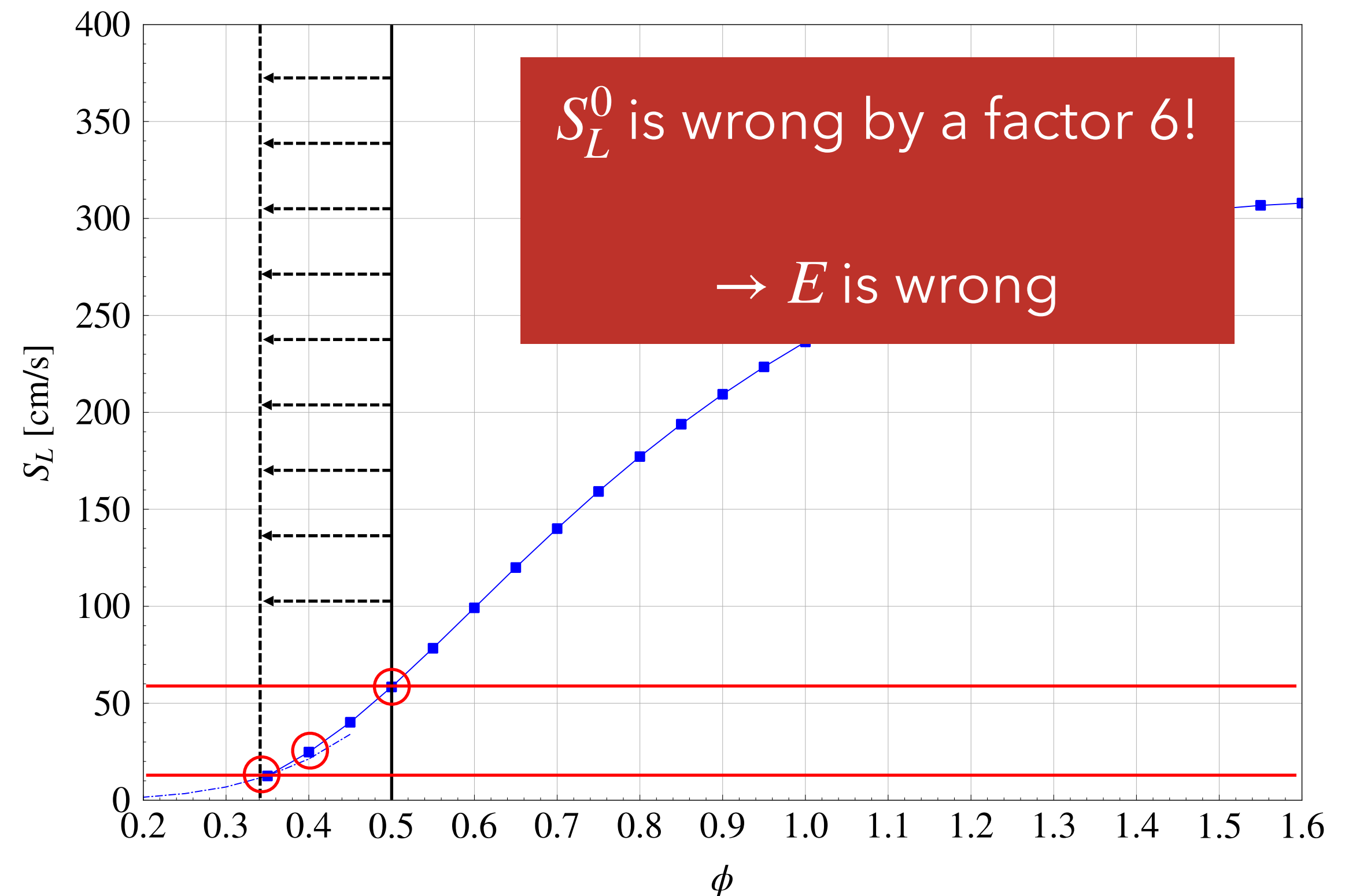
# Thickened Flame model: TFLES

*Equivalence ratio based on Bilger mixture fraction*

## Laminar flame speed



Lean H<sub>2</sub>-Air Flame ( $\phi = 0.5$ ,  $T_u = 300$  K,  $P = 1$  bar)



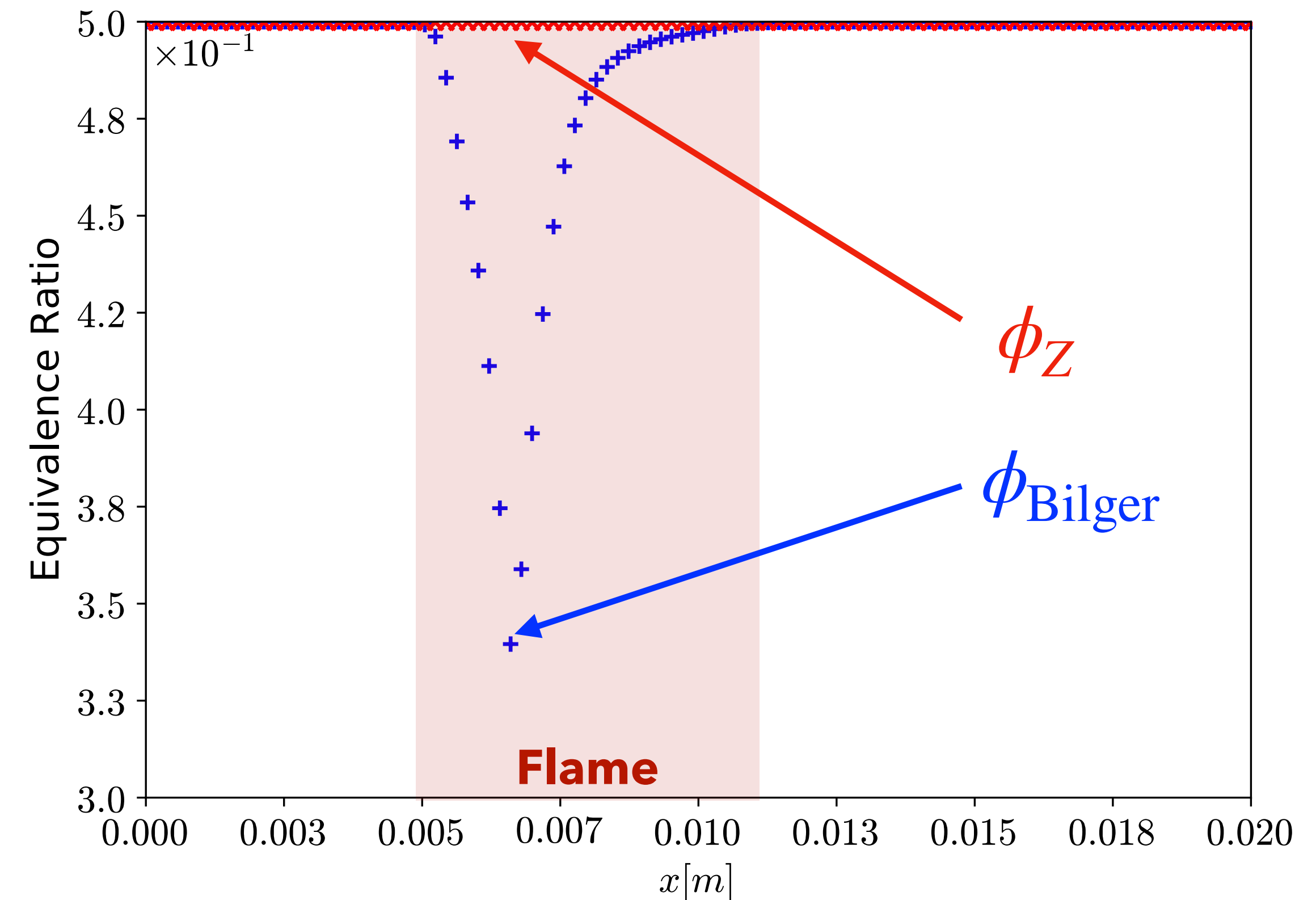
# Thickened Flame model: TFLES

*Solution: equivalence ratio based on a passive scalar*

- Solving an additional transport equation for a passive scalar:

$$\frac{\partial \rho Z}{\partial t} + \frac{\partial}{\partial x_i} (\rho u_i Z) = \frac{\partial}{\partial x_i} \left( \rho D_Z \frac{\partial Z}{\partial x_i} \right)$$

- $Z$  is used (instead of  $Z_{\text{Bilger}}$ ) to recover the equivalence ratio

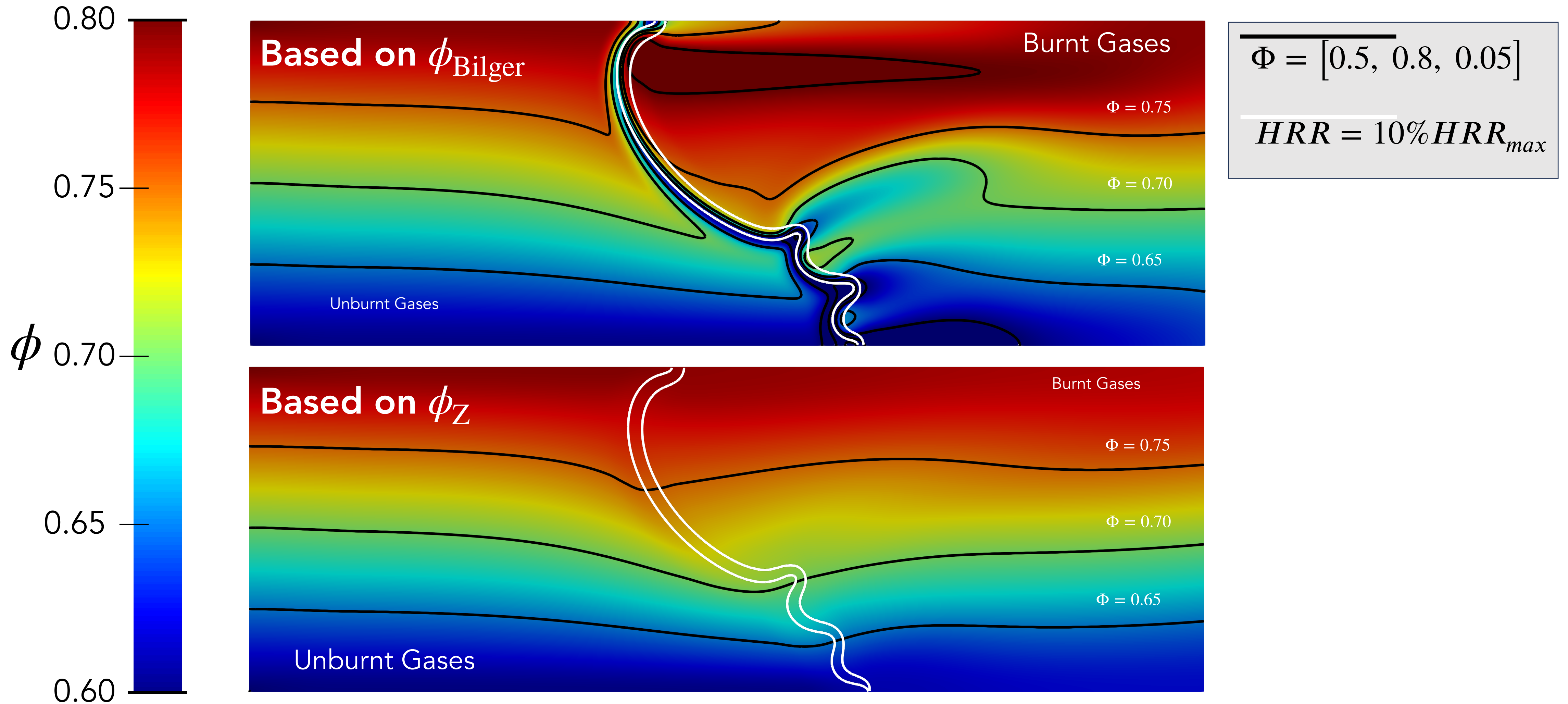




# Thickened Flame model: TFLES

*Solution: equivalence ratio based on a passive scalar*

## 2D stratified flame

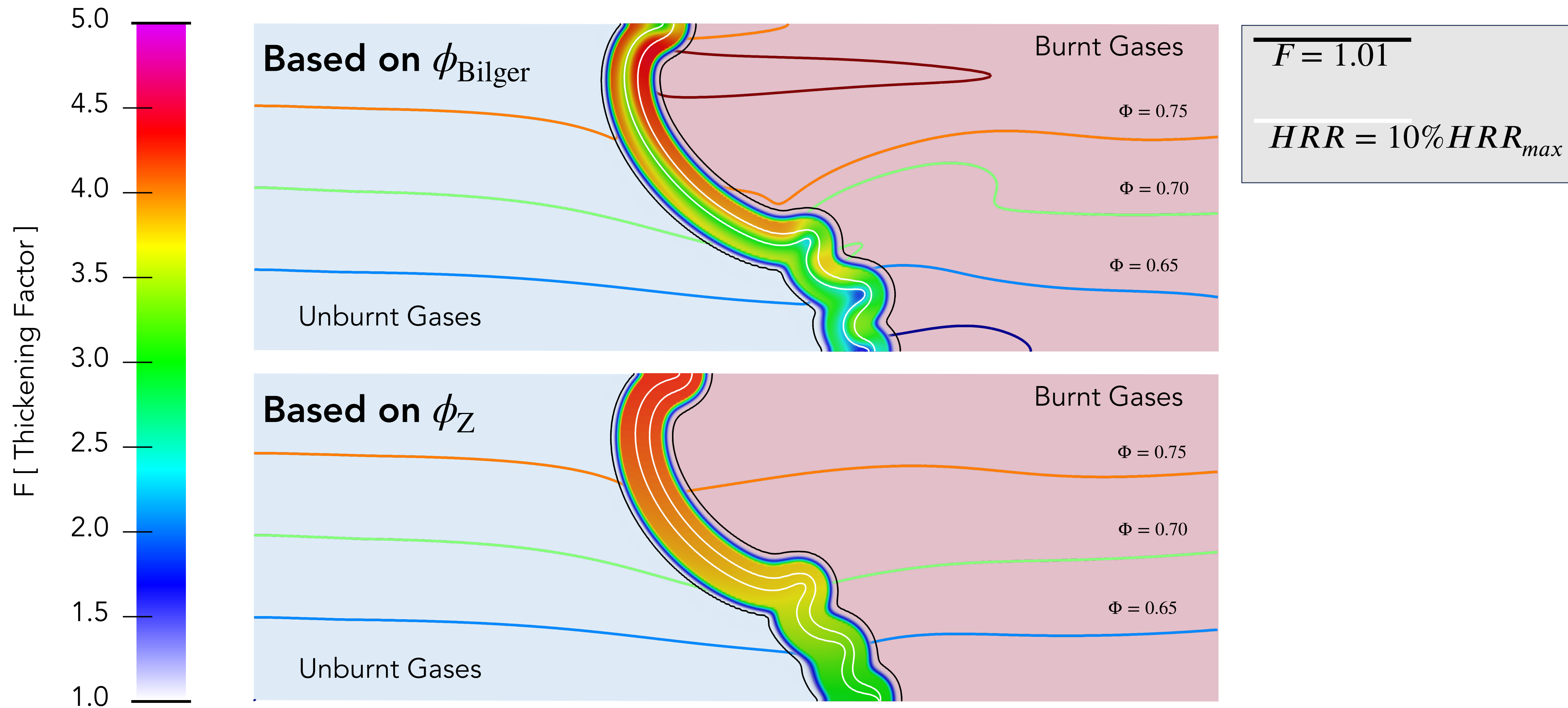




# Thickened Flame model: TFLES

*Solution: equivalence ratio based on a passive scalar*

## 2D stratified flame





# Summary

## *What we discussed today*

### Chemistry

- Contrary to conventional fuel Detailed Chemistry is affordable for  $H_2$

### Transport properties

- Usual viscosity laws are wrong in pure  $H_2$ 
  - Specific correction for  $H_2$ /air: Wilke's Law
    - Variable Prandtl and Schmidt numbers to recover species and heat diffusion properties
- Mixture-averaged transport works for every mixtures  $\rightarrow$  correct  $\mu, D_k, \lambda$

### Non-adiabatic walls

- Inert non-adiabatic walls are ill-posed
- Infinitely Fast Heterogenous Chemistry fixes this issue

### Thickening Flame model

- Thickening factor  $F$  and efficiency function  $E$  are key parameters
- $F, E$  depend on model and physical parameters
  - $T_w, P, \phi$
- Bilger's mixture fraction is inaccurate in flame front
  - Need to add an additional transport equation
  - Accurate values of  $\phi$  can then be recovered in the flame front



# What we do not have time to discuss today

## *Other modeling aspects*

### Stretch response

- Consumption speed of a  $H_2$  flames strongly affected by stretch (strain & curvature)
  - ◉ Classical TFLES does not properly model this
  - ◉ A model for the stretch response has been created for hydrocarbons [1] and  $H_2$  [2]

### Thermodiffusive (TD) instabilities

- Lean laminar  $H_2$  flames are thermodiffusively unstable
  - ◉ A simple model (Aniello []) based on correlation of Berger [] is available in AVBP
  - ◉ WIP: Coupling TD models to corrected stretch response model gives promising results on laminar spherical flames [2]

### Flame/turbulence interaction

- Charlette efficiency function used in most TFLES simulations
  - ◉ Based on the fractal hypothesis
  - ◉ Are  $H_2$ /Air flame fractal?
    - If so, do they have similar fractal properties than other flames?
- What happens when lean flame become turbulent?
- Do turbulence kill the instabilities or are they still present at small scales?
- Are TD effects negligible compared to turbulent wrinkling ?
- Do they couple ?

Next talk!



# Thank you for your attention

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